





ATLA



LHC Searches

LHC 27 km

CERN R

Ludivine Ceard (CP3, UCLouvain)

Petnica Summer School 26th / 27th of July 2014



CERN Mer

Université catholique de Louvain



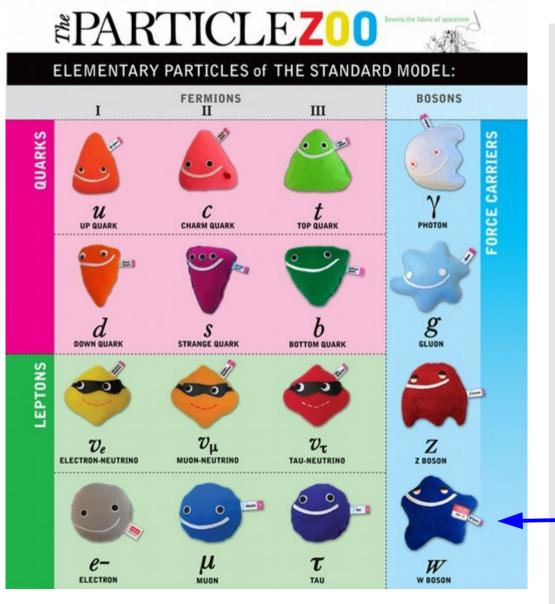
PART I

How did we end up building such a thing as the LHC?

- Accelerators : why?
- Accelerators : What?

• The LHC

- CERN's accelerators complex
- Technology challenges
- Magnets
- The LHC Story
- Lumi and Pile-Up
- Data Taking and Grid
- The Structure of an event
- Analysis of the data in detectors
- What is an Analysis?



- **Standard Model** : seen on your previous course, seems like a good theory but how to verify the existence of the predicted particles?
- Accelerators by accelerating and colliding the particles allows the creation of particles we couldn't observe otherwise because they decay too fast.
- High energy allows to create high mass particles

Most of them observed in the past : USA (Tevatron, SLAC), CERN (SPS + LEP).

State of observation before LHC starts.

• What is left to be answered?



To be answered :

the Higgs prediction : (flashback before 2012) last particle predicted by the SM with the Higgs field, responsible for the mass of all particles.
 But hasn't been observed yet : Does it exist? What are the properties?
 Should we look for another mass explanation?



• Matter-antimatter : there should be in equal amount, but the matter dominates. We have upper limit on amount of anti matter from gamma-rays and cosmic microwave background.

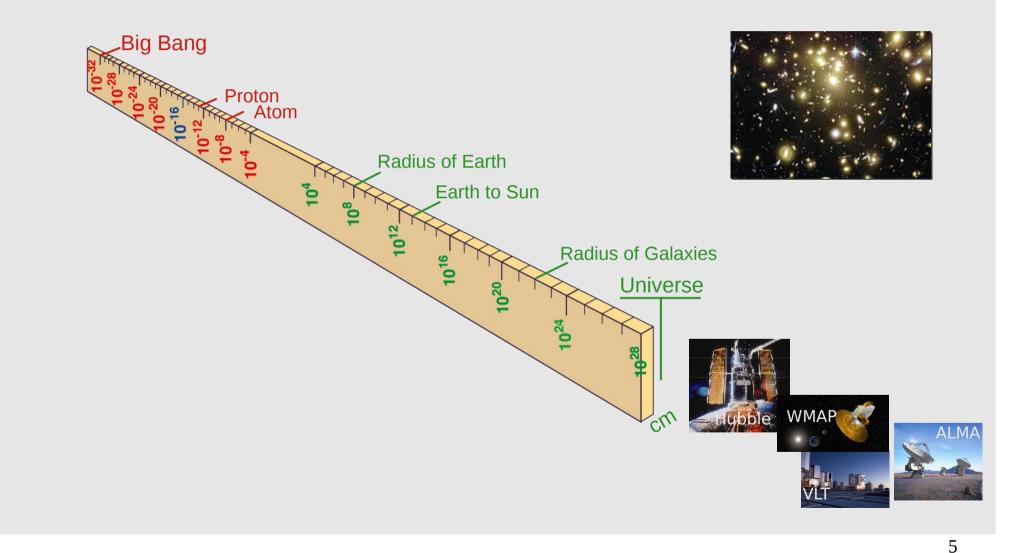
• dark matter : cosmo and astro observations showed that all visible matter ~ 4% of the Universe. Search for particle or explanation responsible for dark matter (23%, supersymmetric particles?) and dark energy (73%).

 force unification : how does gravity fit into the picture? SM does not offer a unified description of all fundamental forces : gravity cannot be described like the 3 others. Supersymmetry (with massive partner for SM particles?)

 Quark-Gluon plasma : a window to the early stage of our universe. After the Big Bang, the Universe went through a stage where matter existed as a hot dense soup of elementary particles. Cooling → quarks trapped (confinement). Can reproduce the QGP by colliding heavy ions. T ~ 1M * Tsun

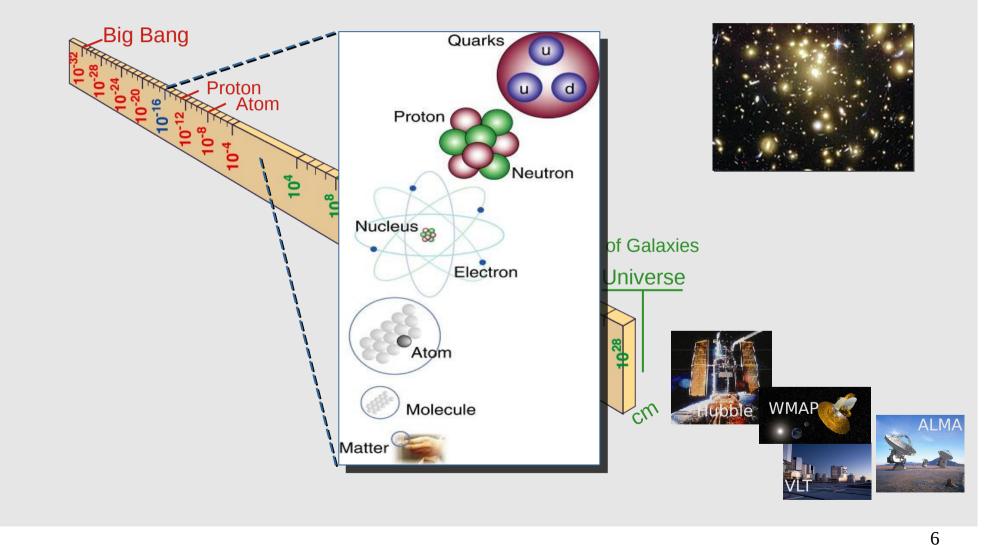


Broadest scale

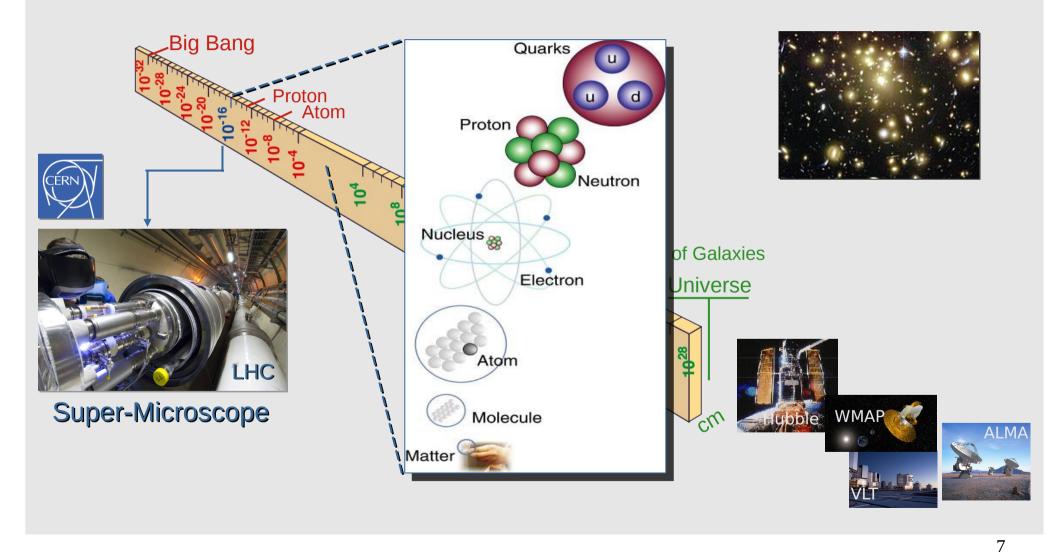




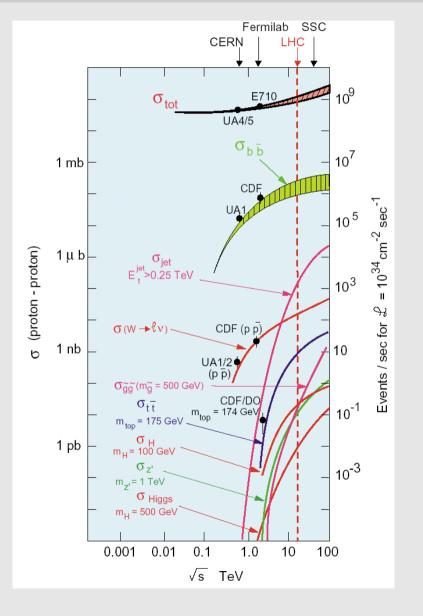
Interesting = rare , heavy = energy



Broadest scale



Interesting = rare , heavy = energy



• The interesting processes are rare = small cross section = small probability to happen.

• The cross sections of the production of those interesting events containing heavy particles increases dramatically with the energy at the center of the collision.

Types :

hadrons -> highest energy but initial state not precisely known due to proton being a composite particle from quarks and gluons

leptons -> for precision physics since initial state of collision known precisely. But energy not so high.



Very first ones :

Electrostatic field do not change with time. Disadvantage = large electric field needed to accelerate particles to experimentally useful E.



Van de Graff 1931 2MeV

Cockcroft-Walton 1932



And later :

Oscillating : require fields that periodically change with the time. → acceleration to extreme high E

Linear Accelerators

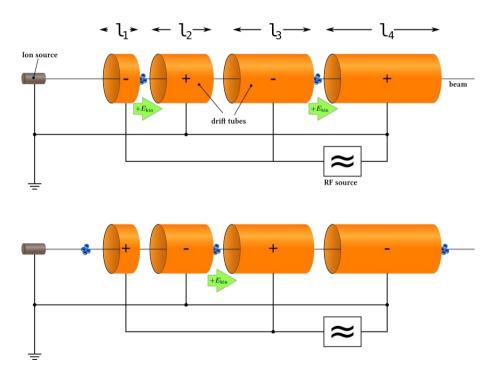
LINAC

Basic principles of LinAcs unchanged since Wilderoe (1920's)

• Alternative current (periodically reversed flow of electric charge) and series of drift tubes.

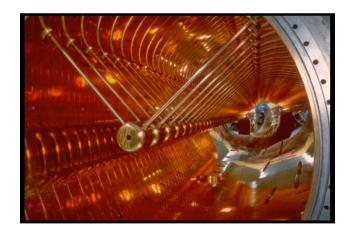
 \rightarrow particle accelerated during a peak of voltage and hidden in drift tube during anti peak to avoid come back to starting point

• As particle get faster the drift tube need to get longer : length is one limiting factor for high energies



Largest LINAC at SLAC US : 3.2km, $e^{\scriptscriptstyle -}$ and $e^{\scriptscriptstyle +}$, 50 GeV

LINAC 4 being build at CERN



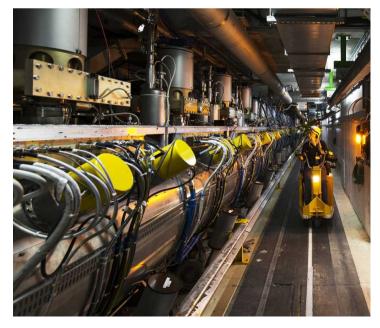
Need to synchronize between particle speed and electrical field. Synchronous phase \rightarrow creation of bunches of particles.

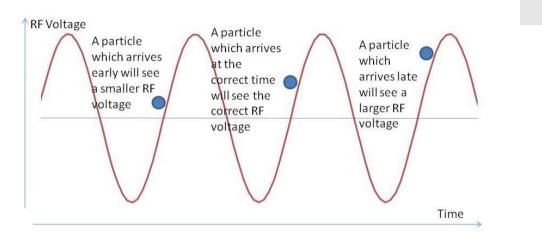
Linear Accelerators

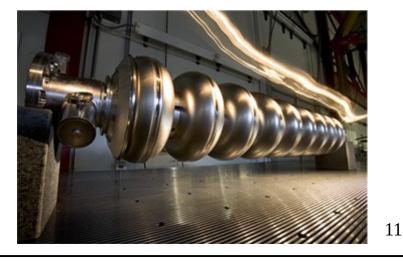
LINAC

As $v_{particle} \sim speed of light the switching rate of the electric fields becomes so high that they operate at radio frequencies and so resonant microwave cavities are used in higher energy machines instead of simple drift tubes.$

RF cavities at LHC



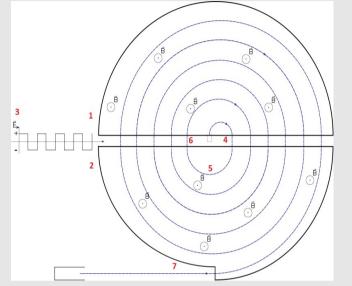






Cyclotron

Cyclotron



First cyclotron: E.Lawrence in the 30's 2 D-shape electrodes alternatively charged by oscillator.

- **Principle** : magnetic field applied perpendicular to the plane of motion of an accelerated particle.
 - \rightarrow even more accelerated
- Constant magnetic field.
- Classical mechanics. 25 GeV max!

Improvement

Take into account the effects of relativity on mass of particles approaching the speed of light .

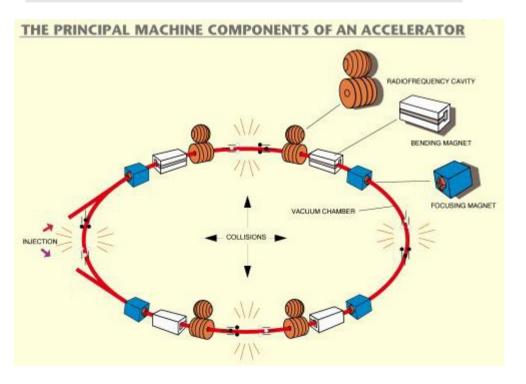
 \rightarrow 1945, called Microtron.



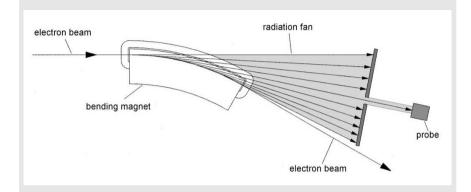
Synchrotron

Synchrotron

- Idea = keep the accelerated particles on a constant orbital radius.
- Synchronize the B with the energy of the accelerated particles.
- First in 1954. Many since!



Synchrotron radiation



• An electromagnetic radiation emitted by an accelerated particle with v ~c and a trajectory bended by a magnetic field.

 Synchrotron radiation : energy loss ~ 1/mass⁴

Synchrotron

Notation \sqrt{s} = blabla eV

Beam energies

1.) reminder of some relativistic formula

total energy $E^2 = p^2 c^2 + m_0^2 c^4$

$$\Rightarrow cp = \sqrt{E^2 - m_0^2 c^4} = \sqrt{(\gamma m_0 c^2)^2 - (m_0 c^2)^2} = \sqrt{\gamma^2 - 1} m_0 c^2$$

 $cp = \gamma \beta * m_0 c^2$

Proton-antiproton collider Tevatron (1980s-2010)

2.) energy balance of colliding particles

rest energy of a particle $E_0^2 = (m_0 c^2)^2 = E^2 - p^2 c^2$

in exactly the same way we define a center of mass energy of a system of particles:

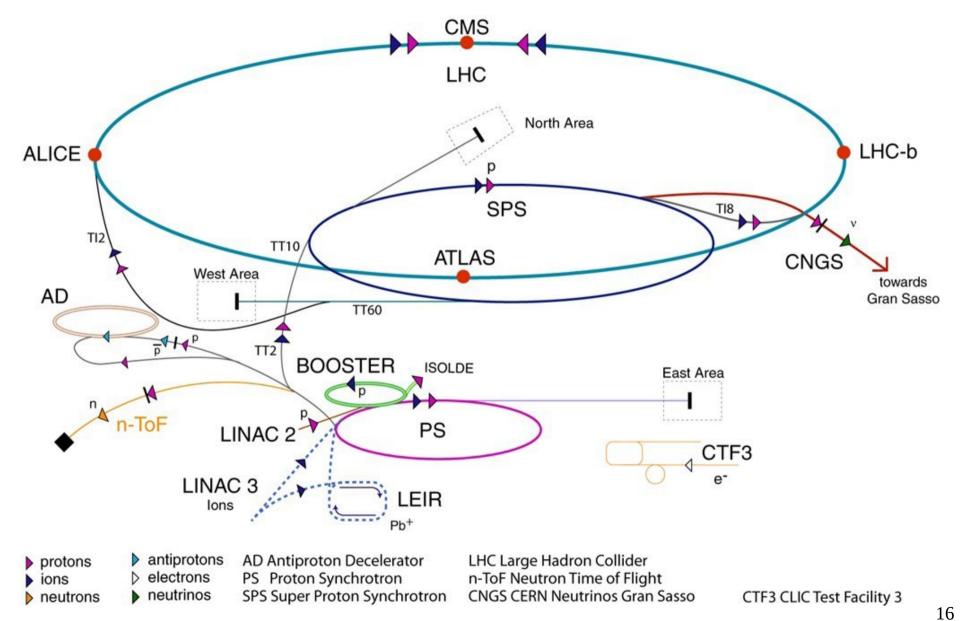
$$E_{cm}^{2} = \left(\sum_{i} E_{i}\right)^{2} - \left(\sum_{i} cp_{i}\right)^{2}$$

Proton-antiproton collider Tevatron (1980s-2010)

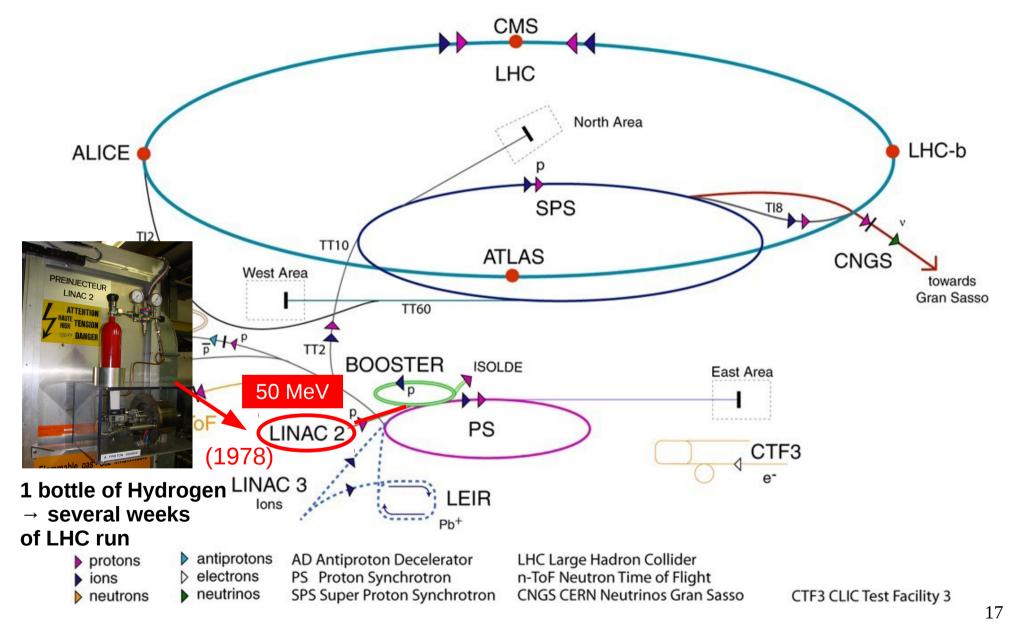
- Creates antiprotons by colliding protons on Ni Target
- Highest energy achieved with particle against anti-particle: 1.96 TeV
- But difficulties to produce and store them
 Solution =
 proton against proton



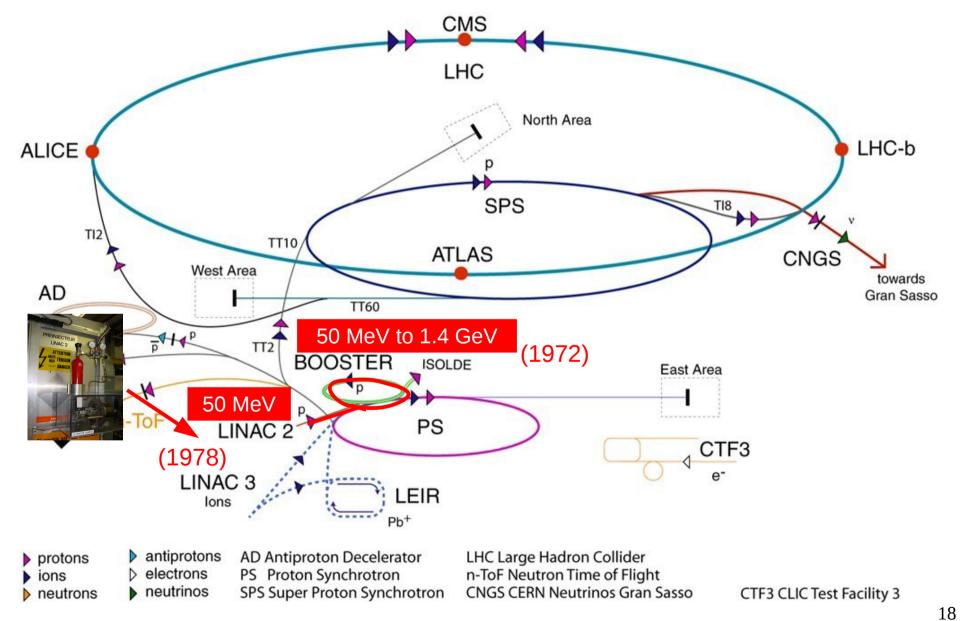




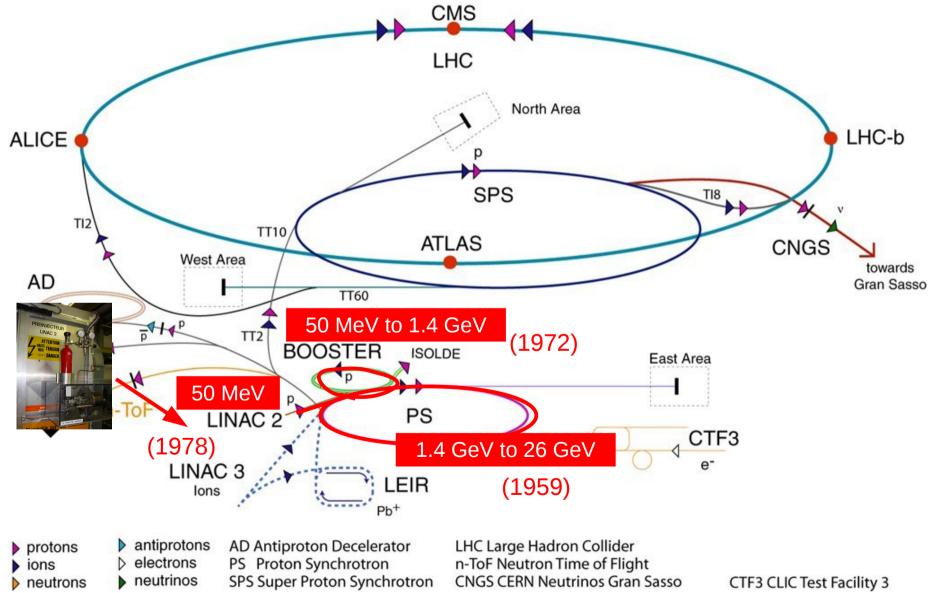






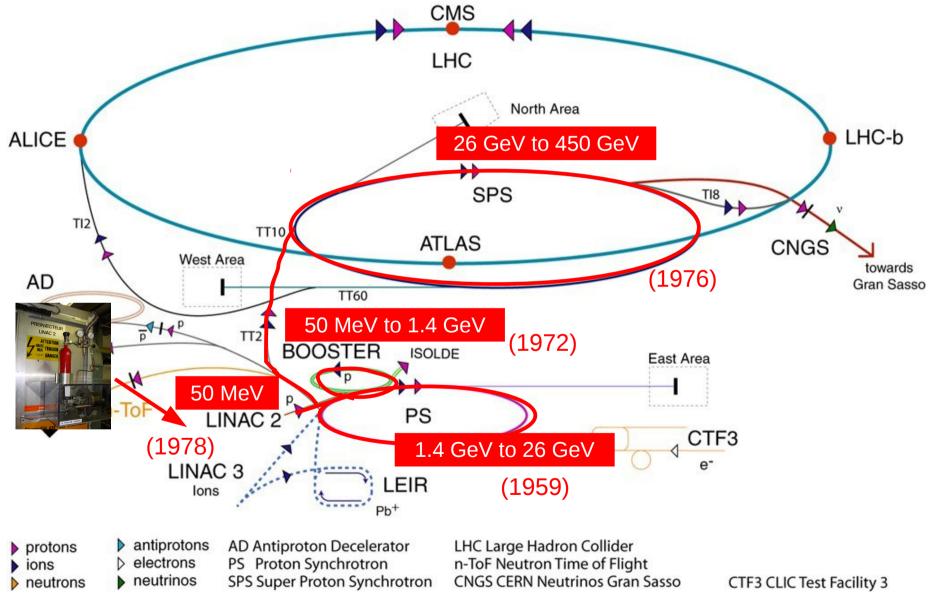






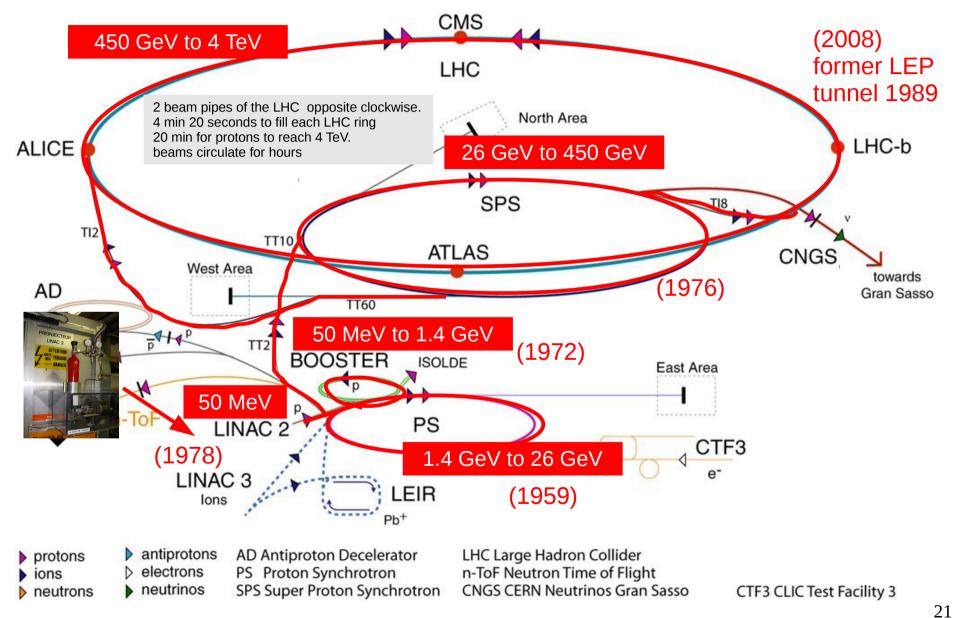
19





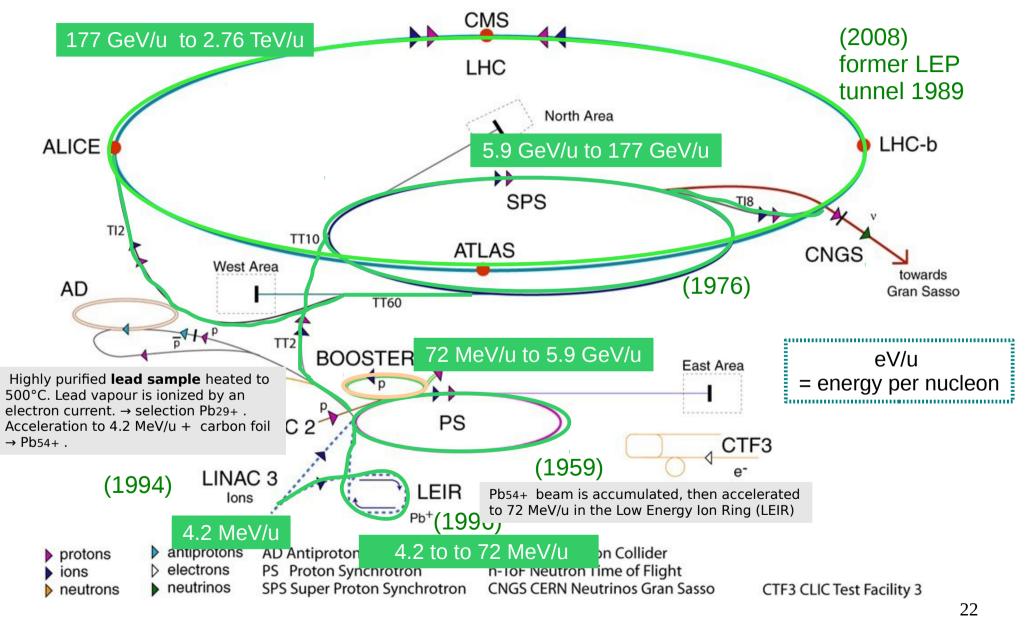
20





lons acceleration







Technology challenges



Numbers!

Several thousand billion protons 99.9999991% of light speed Orbit 27km ring 1 turn : 11 000 times/second A billion collisions a second

The beam vacuum pressure : 10⁻¹³ atm, to avoid collisions with gas molecules : vacuum similar to interplanetary space: pressure in the beam-pipes will be ten times lower than on the Moon

Beams squeezed to about 16 μ m (a human hair is about 50 μ m thick).

LHC 1.9 degrees above absolute zero = - 271 C Outer space 2.7 degrees above zero = - 270 C

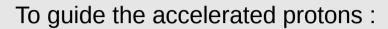
The total energy in each beam at maximum energy is about 350 MJ, which is about as energetic as a 400 tonnes train, like the French TGV, traveling at 150 km/h. This is enough energy to melt around 500 kg of copper.

23

Technology challenge



Dipole magnets



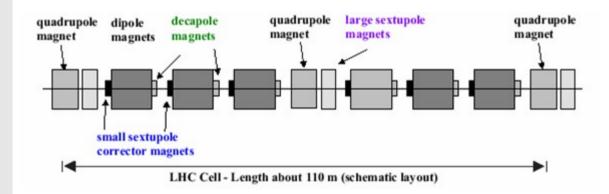
```
7000 GeV Proton storage ring
dipole magnets N = 1232
I = 15 \text{ m}
q = +1 \text{ e}
Example LHC:
```

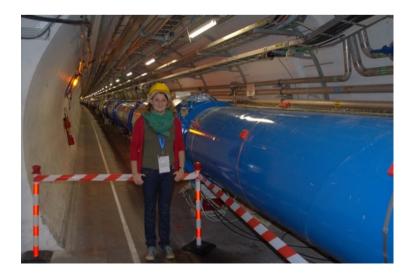
$$\int B \, dl \approx N \, l \, B = 2\pi \, p/e$$

$$B \approx \frac{2\pi \, 7000 \, 10^9 \, eV}{1232 \, 15 \, m \, 3 \, 10^8 \, \frac{m}{s} \, e} = 8.3 \, Teslo$$

more than 100,000 times
 more powerful than the Earth's
 magnetic field!

Designed for 14 TeV.





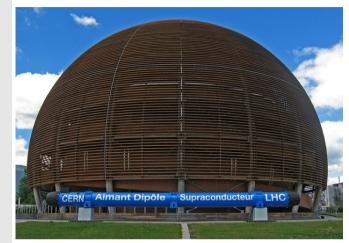
24

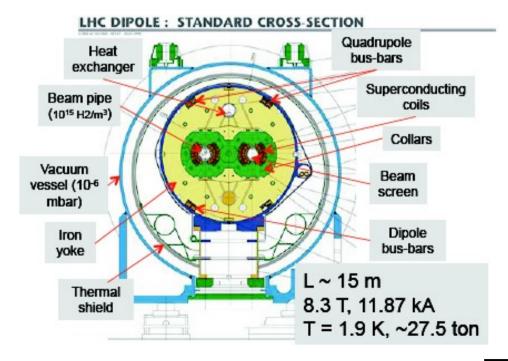
More on Magnets

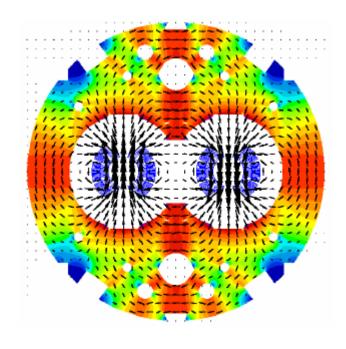


How do they do?

- LHC dipoles use niobium-titanium (NbTi) cables, which become superconducting below a temperature of 10 K (–263.2°C)
 → conduct electricity without resistance.
- LHC operate at 1.9 K
- \bullet Current of 11 850 A flows in the dipoles, to create the high magnetic field of 8.33 T







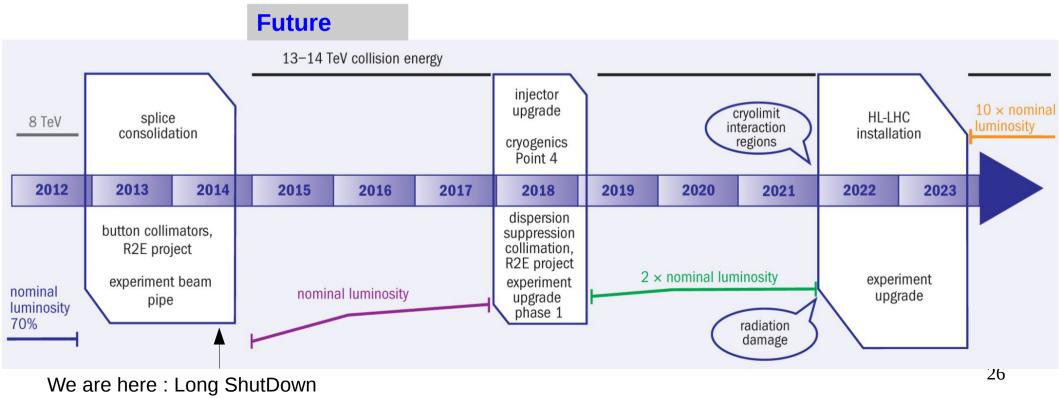


The LHC story



Past

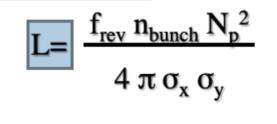
10 September 2008	: first beam!		
19 September 2008	: LHC accident , faulty elect	strical bus connection \rightarrow release of He	→ damages
November 2009	: 1.18 TeV per beam		
30 March 2010	: first collisions! at 7 TeV!		
Ap 2010 – Nov 2011	: p-p coll. at 7TeV	Nov 2011-Dec 2011 : Pb-Pb coll.	at 2.76 TeV/u
Ap 2012 -Dec 2012	: p-p coll. at 8TeV	Jan 2013- Feb 2013 : p-Pb coll.	at 5.02 TeV/u



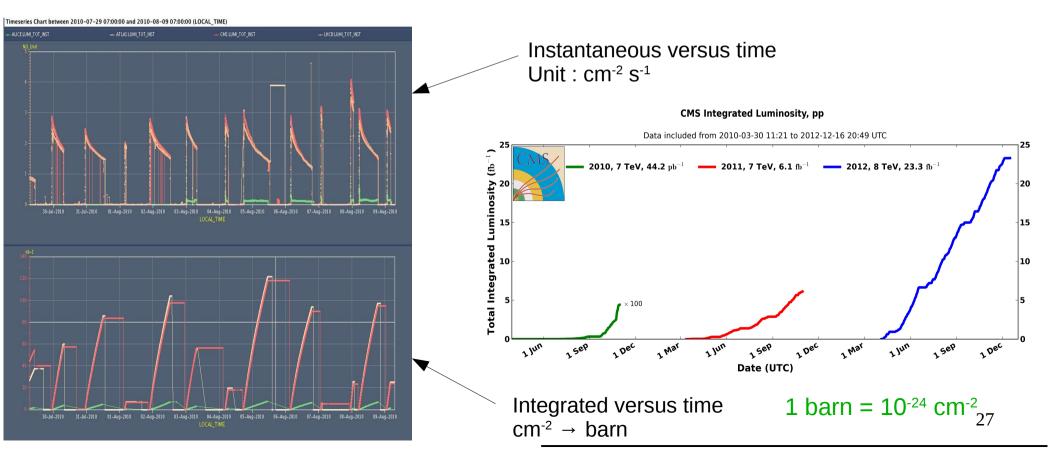


Luminosity



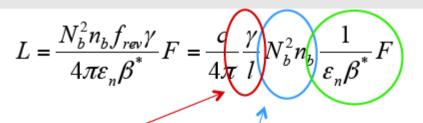


revolving frequency: f_{rev} =11245.5/s #bunches: n_{bunch} =2808 #protons / bunch: N_p = 1.15 x 10¹¹ Area of beams: $4\pi\sigma_x\sigma_v$ ~40 µm



Luminosity





Accelerator features

Energy of the machine 7 TeV Length of the machine 27 km

Beam intensity features

N_b Number of particles per bunch 1.15×10¹¹ n_b Number of bunches ~2808

Nominal luminosity: 10³⁴ cm⁻² s⁻¹

(considered very challenging in the 90's, pushed up to compete with SuperconductingSuperCollider)

Beam geometry features

 ϵ_n Size of the beam from injectors: 3.75 mm mrad β^* Squeeze of the beam in IP (LHC optics): 55 cm F: geometry reduction factor: 0.84



Luminosity



$$L = \frac{N_b^2 n_b f_{rev} \gamma}{4\pi\varepsilon_n \beta^*} F = \frac{q}{4\pi} \frac{\gamma}{l} N_b^2 n_b \frac{1}{\varepsilon_n \beta^*} F$$

			Nominal	September 2012	Gain L	Ultimate	Gain L
_	N_b	(adim)	1.15E+11	1.55E+11	1.82	1.7E+11	2.2
	ε_n	(m rad)	3.75E-06	2.5E-06	1.50	3.75E-06	1.0
	n _b	(adim)	2808	1380	0.49	2808	1.0
	eta^*	(m)	0.55	0.60	0.92	0.55	1.0
_	E	(TeV)	7.0	4.0	0.57	7.0	1.0
	L	$(cm^{-2} s^{-1})$	1.0E+34	7.0E+33	0.70	2.2E+34	2.2
ł	Pile-up'	* (adim)	25	36		55	
	* 80 mbarn cross section assumed						



Pile-Up

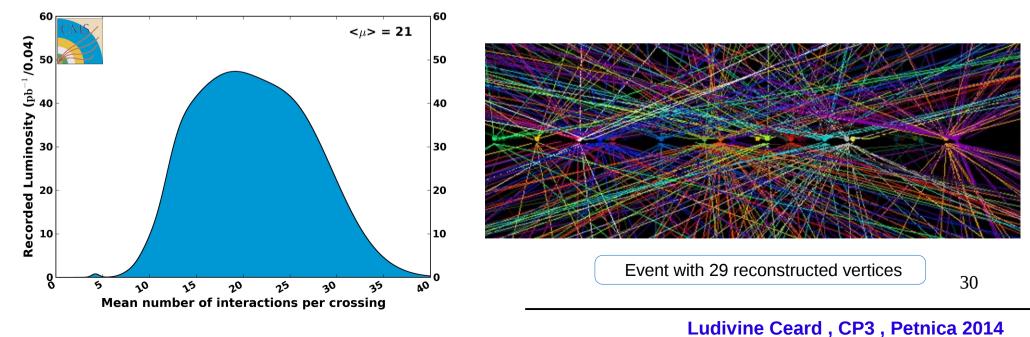


Price to pay for luminosity!

Pile Up is the number of pp collisions in 1 bunch crossing. Characterized by several primary vertices (hard process spot).

Facility	\sqrt{s} [TeV]	$\sigma_{\rm inel} \ [{\rm mb}]$	$L \ [10^{34} \ \mathrm{cm}^{-2} \mathrm{s}^{-1} \]$	τ_b [ns]	$< N_p >$	$< N_p > +2\sigma$
LHC (2012)	8	71.5	.75	50	27	38
LHC (nominal)	14	76	1	25	19	28
LHC (50 ns)	14	76	1	50	38	50

CMS Average Pileup, pp, 2012, $\sqrt{s}=$ 8 TeV





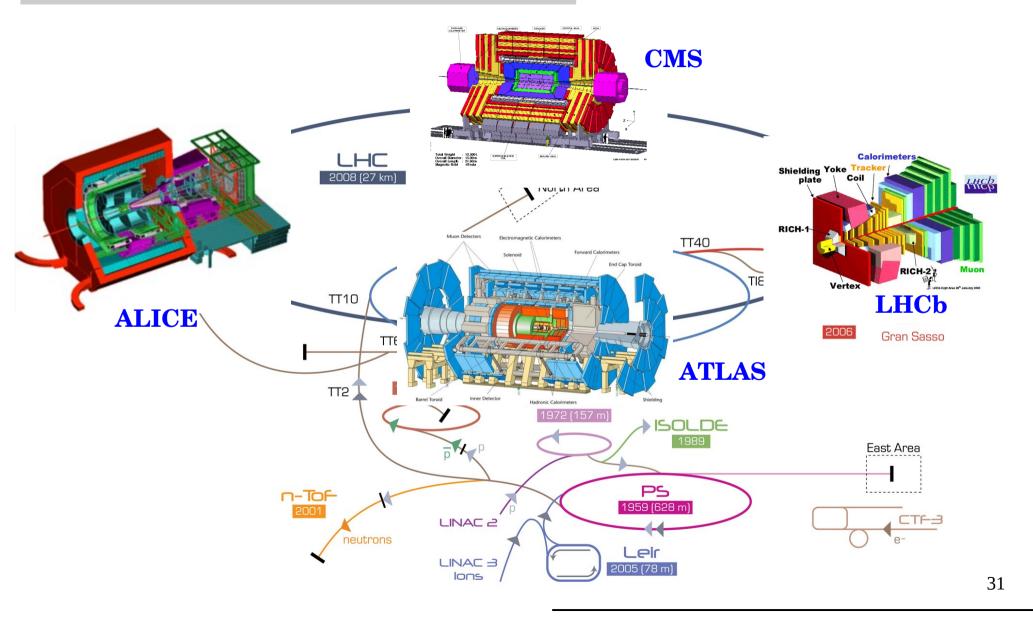


Detectors

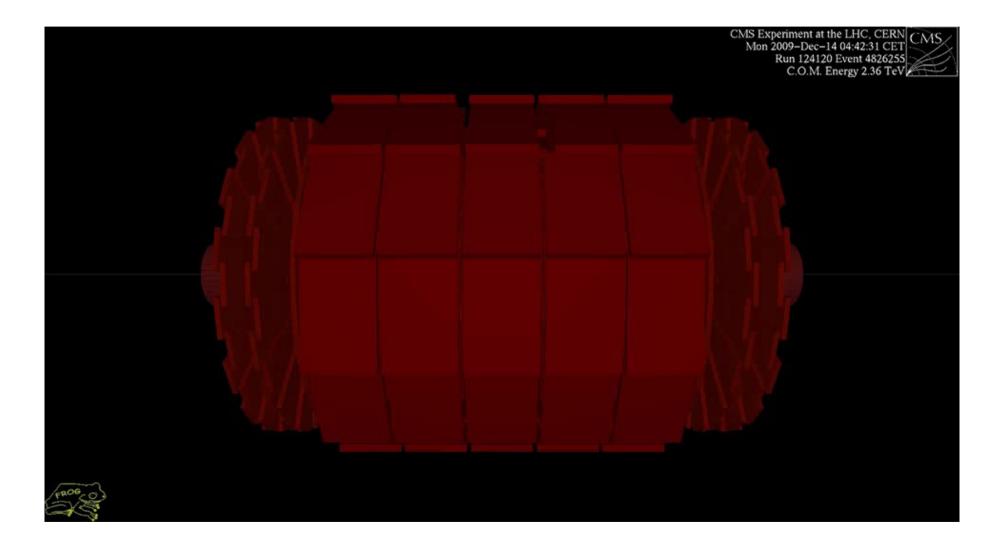




The Four Main experiments of the LHC

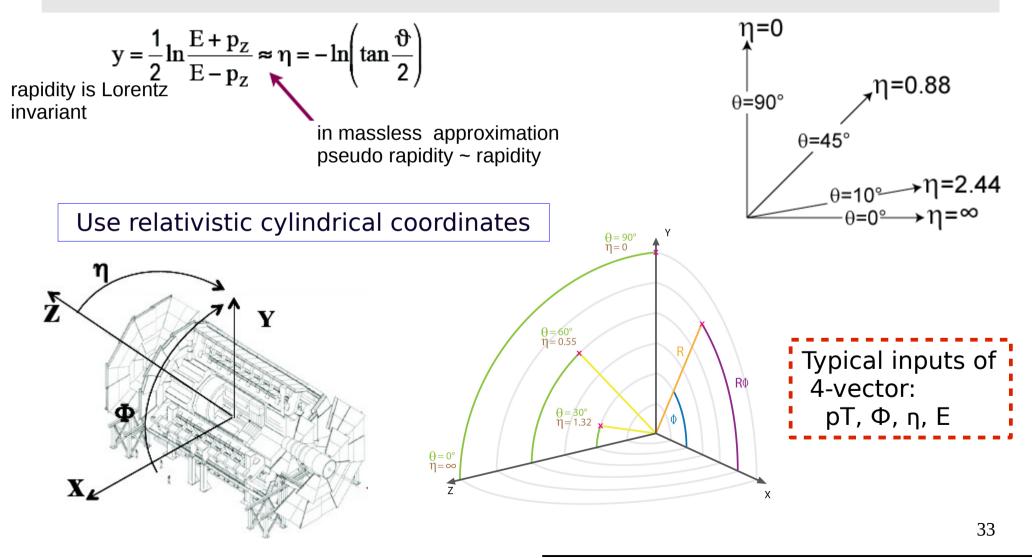


A collision in CMS





As a consequence of the collision kinematic, the visible pZ is not known. Only the conservation of the transverse momentum pT can be used.



Data taking



The LHC computing GRID

 Number of collisions : more than 500 Millions per second ~ 20 MHz

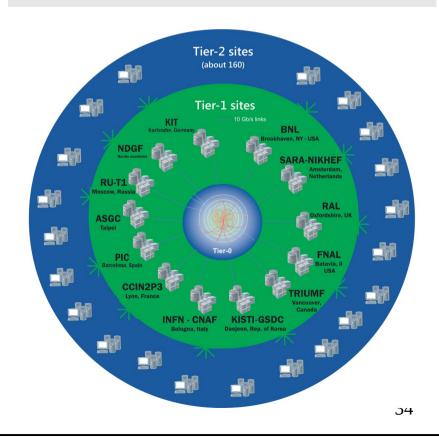
• The LHC experiments represent about 150 million sensors delivering data 40 million times per second.

- After filtering there will be about 100 collisions of interest per second.
- The data flow from all four experiments will be about 700 MB/s ~15 PB per year
- ~ a stack of CDs about 20 km tall each year.

• Enormous amount of data will be accessed and analyzed by thousands of scientists around the world.

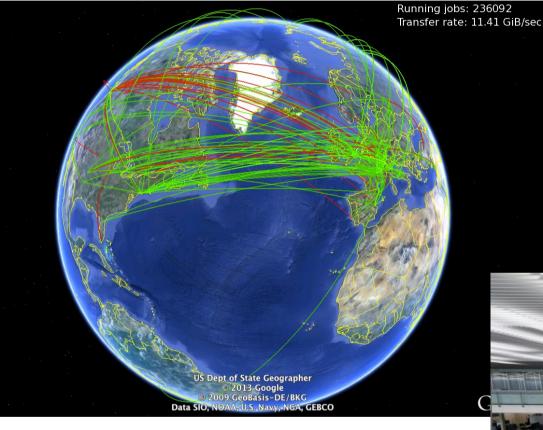
• LHC Computing Grid built and maintain a data storage and analysis infrastructure for the entire high-energy physics community that will use the LHC. 2 levels of Trigger for each exp : • Low level trigger 100 kHz in the cavern, 20 μs to decide → storage underground

- HLT 600 Hz, 100 ms to decide
- \rightarrow storage on the ground \rightarrow T0



The Grid

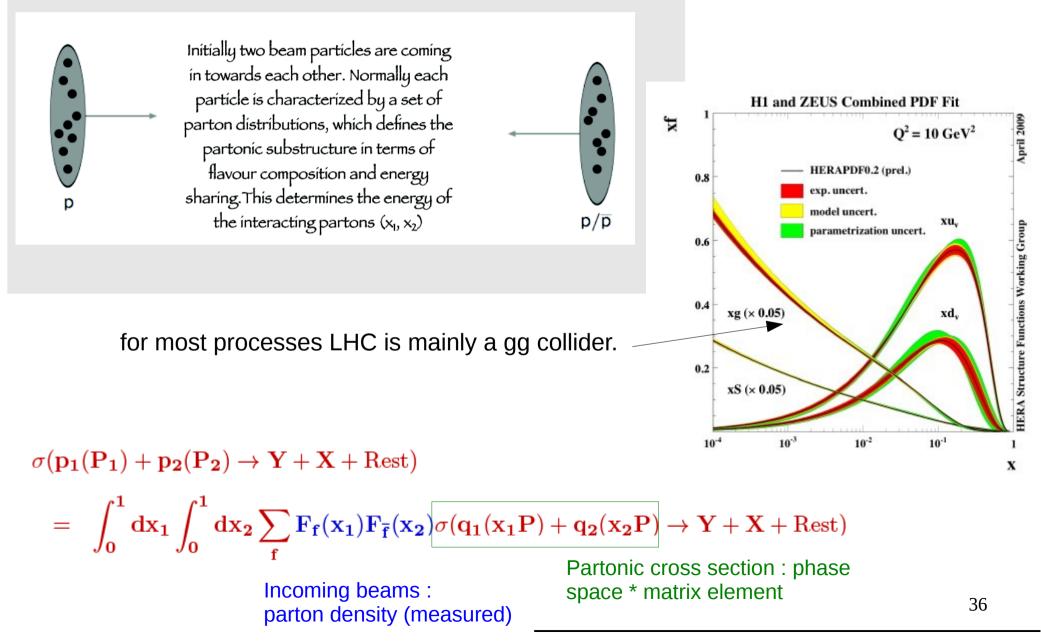




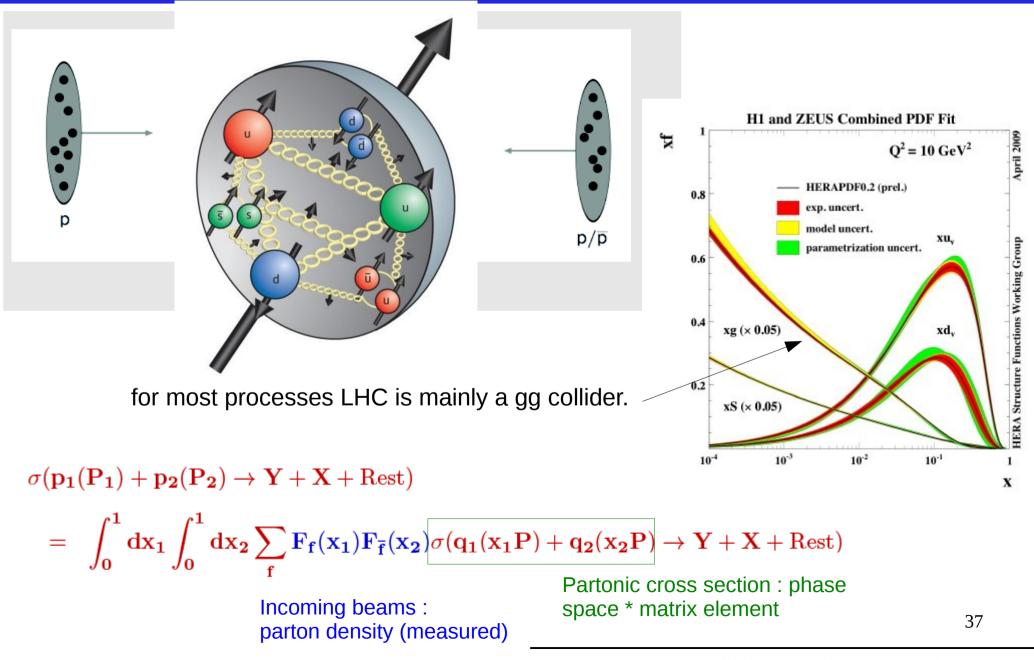
By 2012 data from over 3*10¹⁴ LHC proton-proton collisions had been analyzed. LHC collision data was being produced at approximately 25 petabytes per year, and the LHC Computing Grid had become the world's largest computing grid comprising over 170 computing facilities in a worldwide network across 36 countries



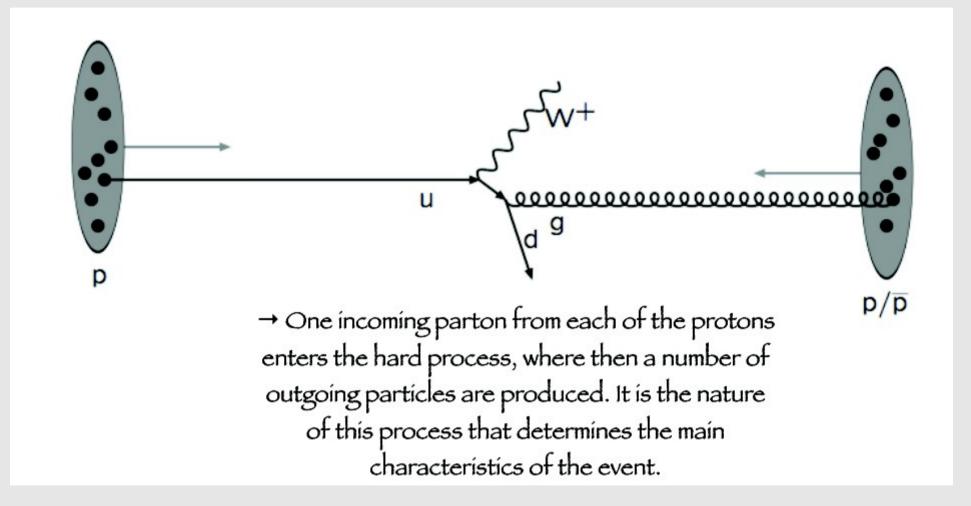
The structure of an event : pdfs



The structure of an event : pdfs

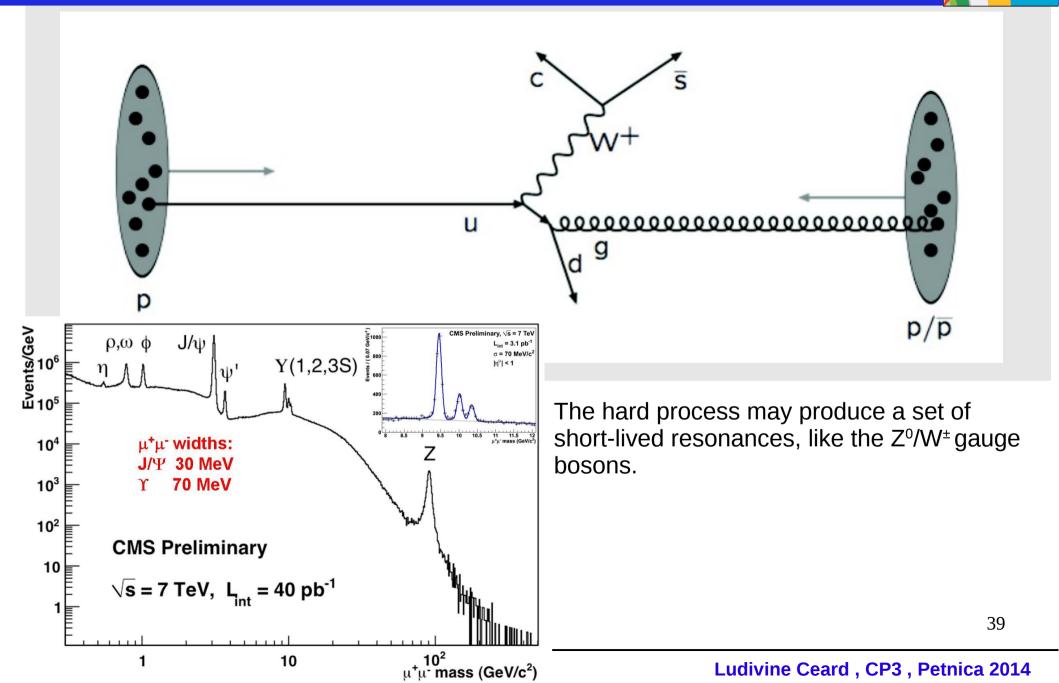


The structure of an event : hard process

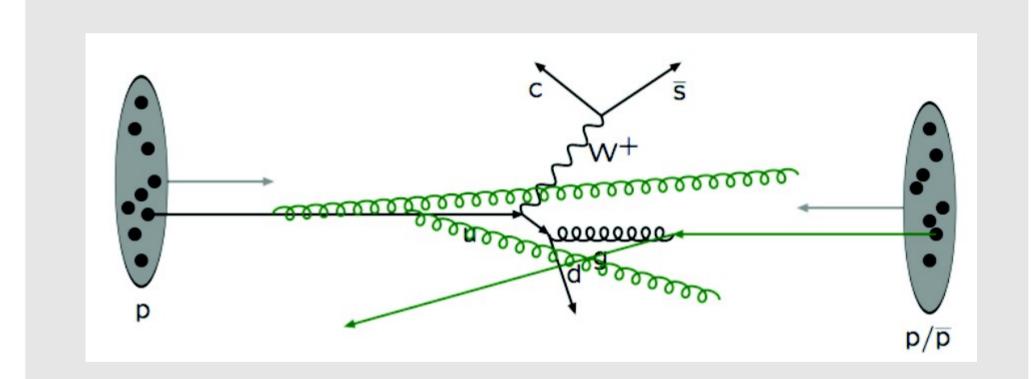


Hard subprocess described by Matrix Elements

The structure of an event : resonance



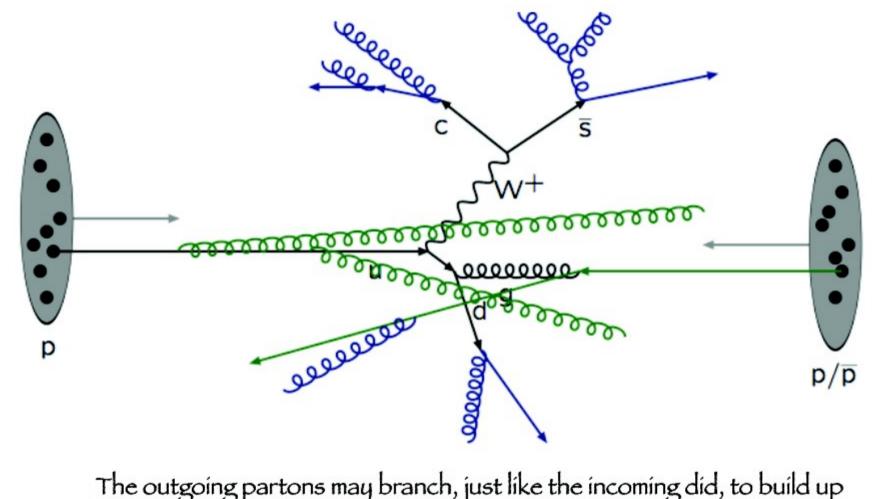




One shower initiator parton from each beam may start off a sequence of branchings, such as $q \rightarrow qg$, which build up an initial-state shower.

Initial State Radiation

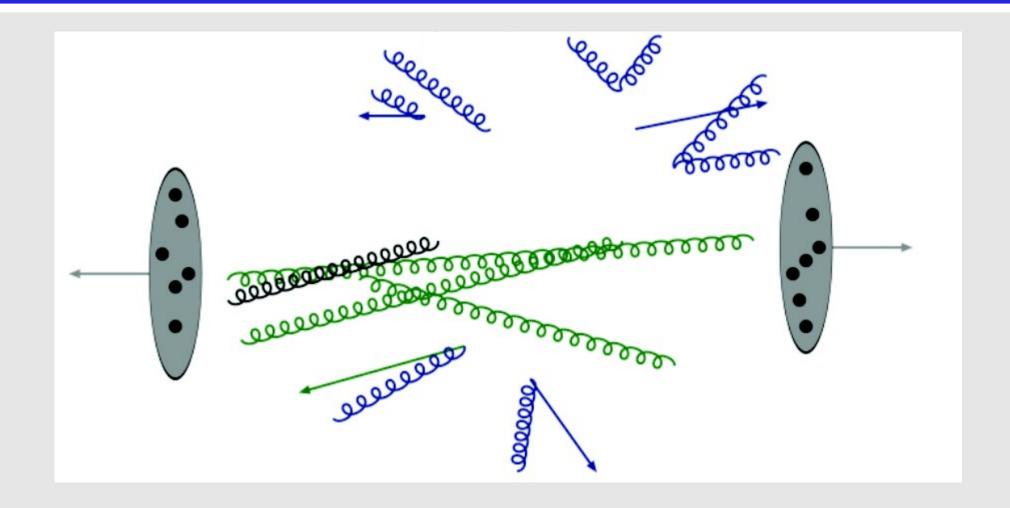
The structure of an event : FSR



The outgoing partons may branch, just like the incoming did, to build up final-state showers.

Final State radiations

An event : underlying event - minimum bias

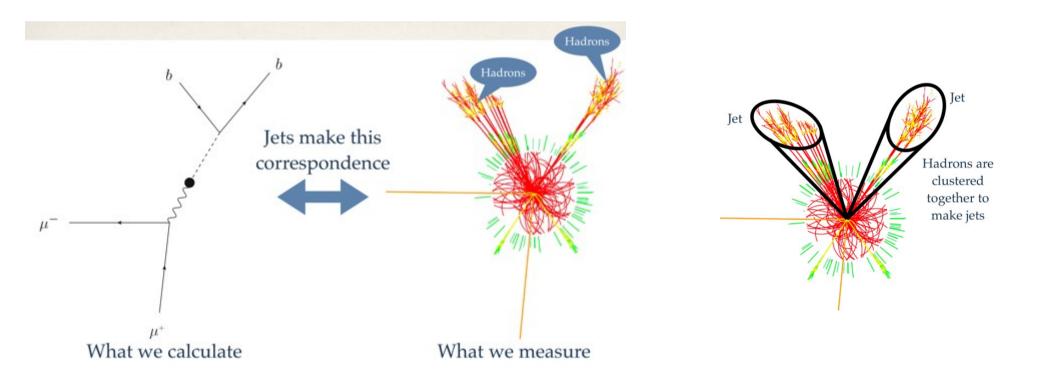


- \bullet Proton remnants (in most cases coloured!) interact: Underlying event, consist of low p_T objects.
- There are events without a hard collision (dependent on pT cutoff) , those are called minimum bias

An event : hadronization

Jets :

Quarks and gluon in the detector cannot be observed as free particles but as many Hadrons and in the detector as a jet of particle in a narrow cone.

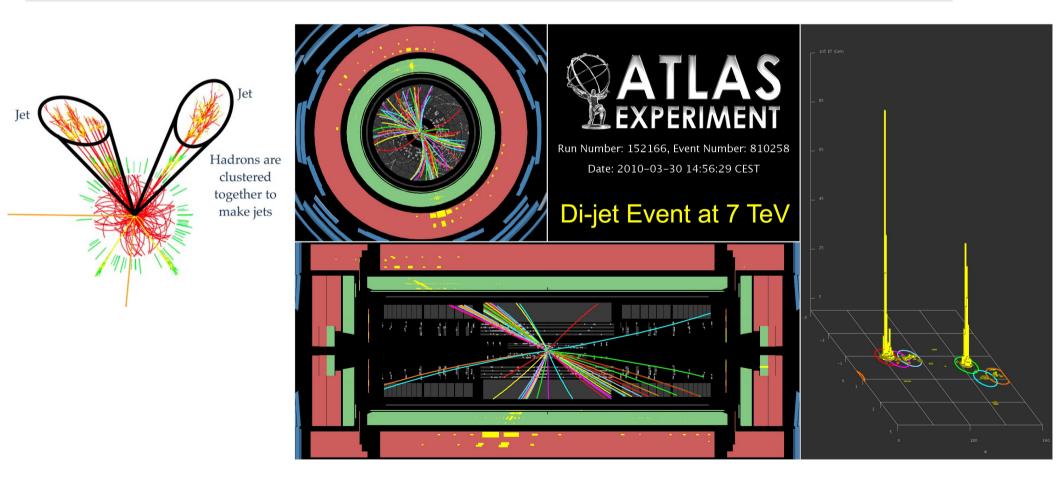


An event : hadronization



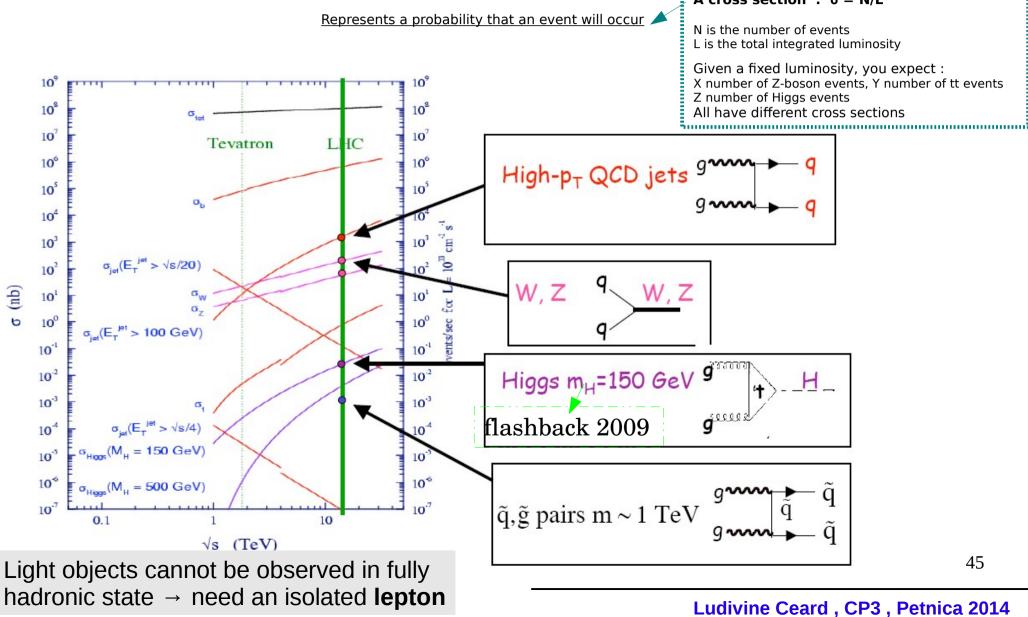
Jets :

Quarks and gluon in the detector cannot be observed as free particles but as many Hadrons and in the detector as a jet of particles in a narrow cone.



QCD and consequences

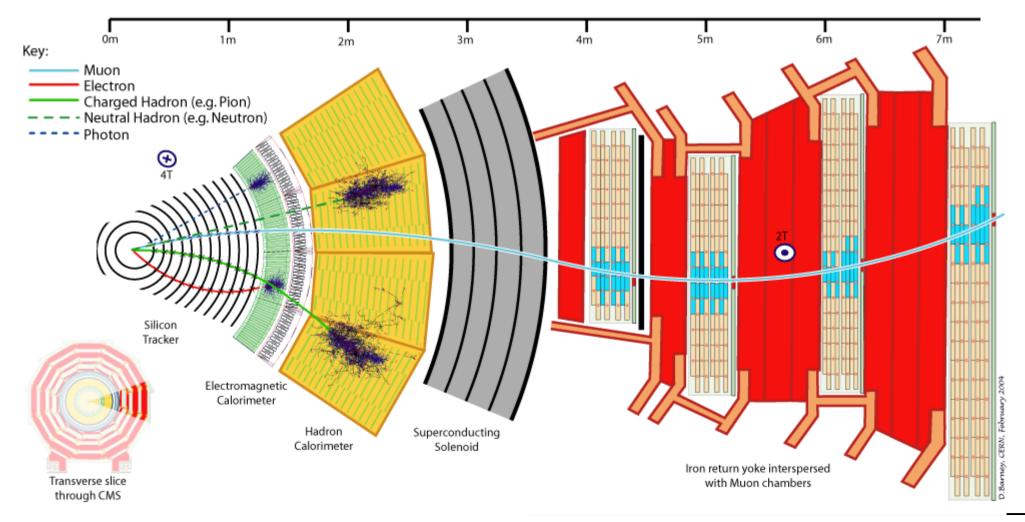
Enormous QCD background compared to what is interesting with jets : A cross section : g = N/L



Analysis of the data

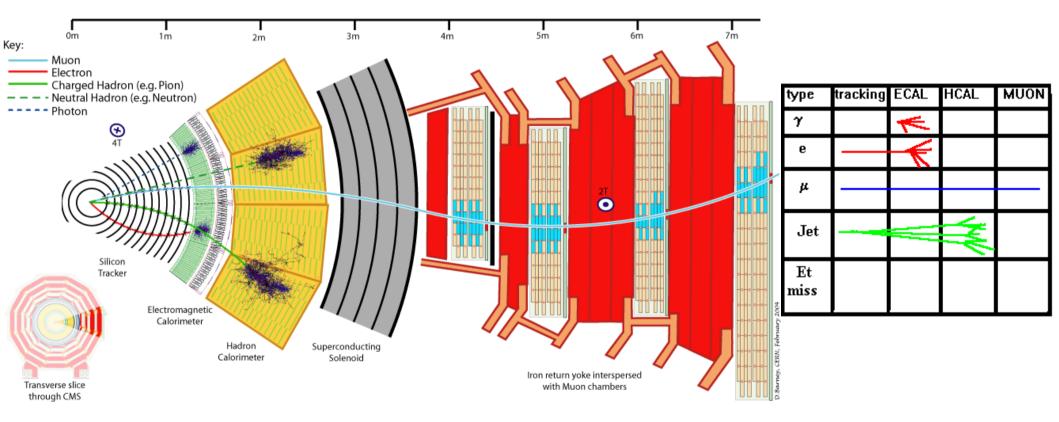
The Final State investigation

- \bullet We can only observe the final state of the particles in the detectors : e, $\mu, \, \gamma$ and hadrons
- We can only reconstruct those final state objects from theirs signatures/tracks left in the detectors



Analysis of the data

The Final State investigation

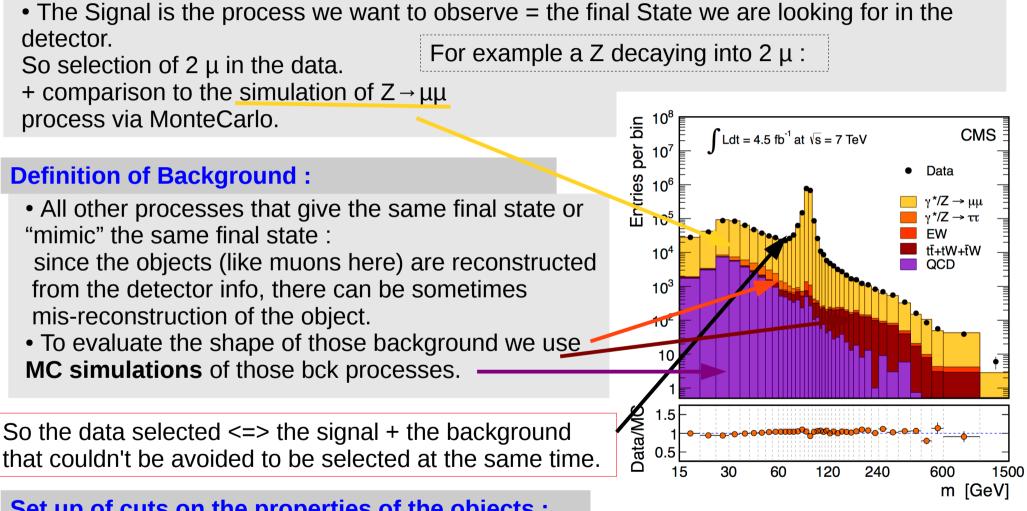


MET = the negative the vector sum of all the transverse components of observed energy. It Indicates the presence of weakly interacting particles, usually neutrinos, but possibly new exotic objects that interact only weakly.

An analysis



Definition of Signal :

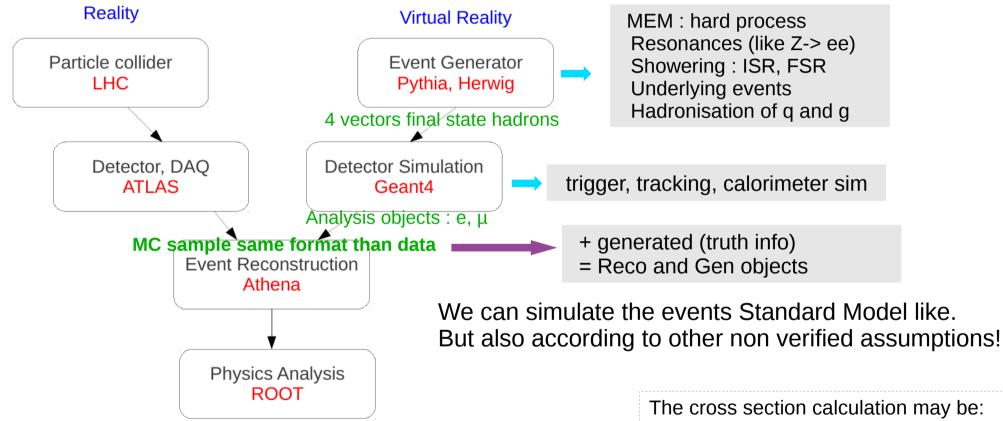


Set up of cuts on the properties of the objects :

Cuts on object properties (pT, Φ , η , E , MET) are decided to keep max of signal while removing max of background.

Monte Carlo Simulation

Simulation of collisions between particles, used in all p.physics experiments :



MC is not the truth !

What happened in the MC generator didn't happen at the LHC The MC is our best guess

The MC only simulates specific physics processes

 \rightarrow when possible comparison of the data with different type of generators

The cross section calculation may be: Wrong, Incomplete, Inaccurate The MC code with unknown bugs in it Higher order corrections which have not been calculated Cannot account for all detector effects

Ludivine Ceard, CP3, Petnica 2014

49



An analysis

Which quantities we are looking at :

Standard Model Physics :

• check of well known quantities and comparison with the MC to get confidence in our measurement (mass of Z for eg)

- check of kinematics variables : pT of objects, combined mass of system (several jets together), MET, Number of jets in the event, angular distribution
- calculation of the cross sections

Searches :

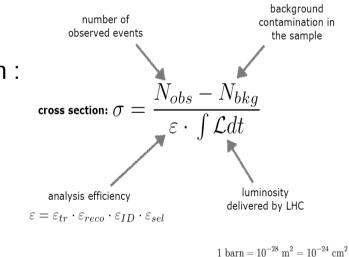
 quantities where some deviation from the SM prediction is expected (thank you theorists) = we look at the data and at the Standard Model MC simulation : can be a prove of new physics.

• Invariant mass for possible unknown resonances

See Higgs example!!

Errors:

Statistical, systematical Efficiencies, purity ...



the cross-section



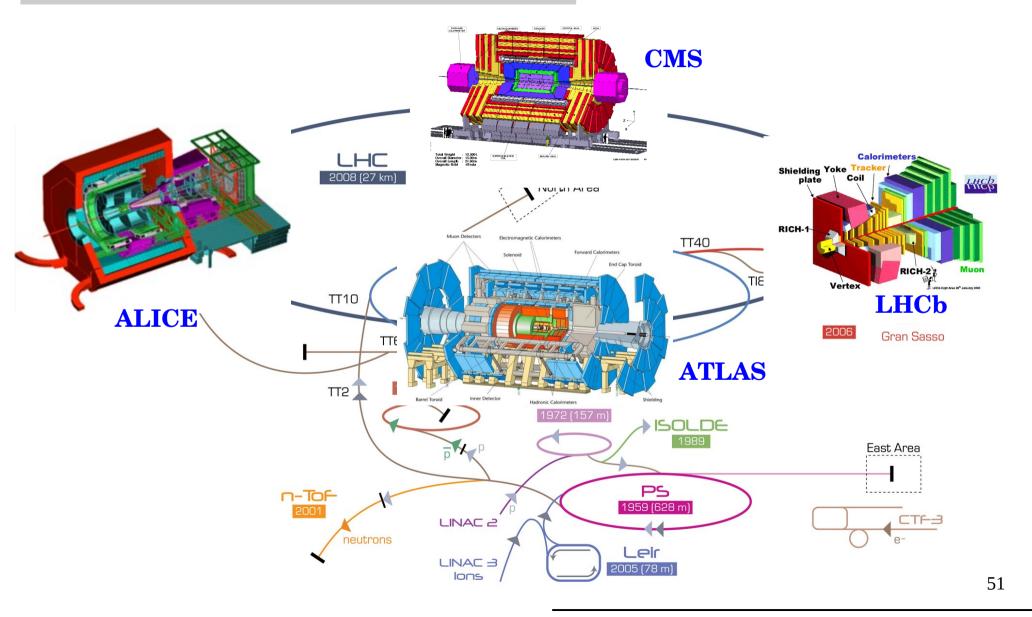


Detectors





The Four Main experiments of the LHC





PART II

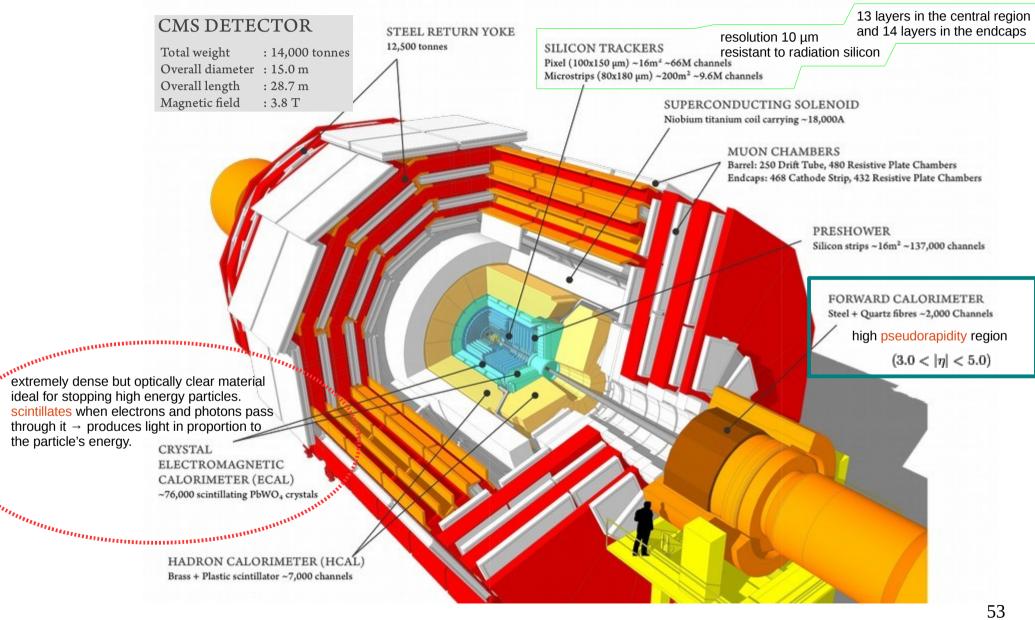
Let's see what people are really doing as Physics at the LHC.

• CMS : Compact Muon Solenoid

- A SMP Analysis from head to toe : the Zbb analysis example
- ATLAS : A Toroidal LHC ApparatuS
 - Search and found for the Higgs into 4 leptons
- The Higgs discovery
- Exotic and SUSY searches
 - Exotic : the example of the Z' particle search
 - Status of Exotica and Susy (for experimentalists)
- LHCb
 - \bullet The $B_{S \ \rightarrow \ \mu\mu}$ result
- ALICE

CMS : Compact Muon Solenoid







Data

VH(125 GeV)

H(125 GeV)

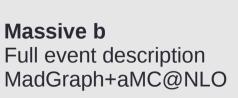
Z + udsco

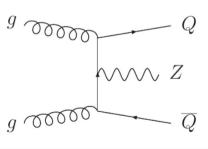
+ bb

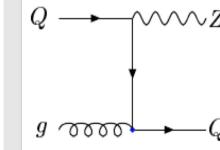
W + udscg Single top

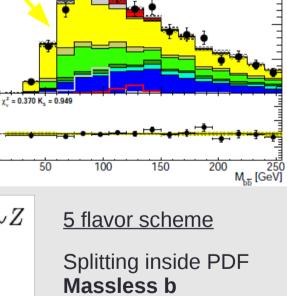
Motivations: why Zbb?

- Z+b-jets is background for many searches for undiscovered processes:
- Events / 15. SM Scalar: Z(II) H(bb) * 300 s = 7 TeV. L = 5.0 fb¹ = 8 TeV. I = 12.1 fb⁻¹ $H \rightarrow Z(II) Z(bb)$ $H \rightarrow Z(II) A(bb)$ **BSM Scalar:** 2HDM/Susy-like * 200 150 ➔ Understand Z+b-jets process **Study kinematics** 100 50 2 1.5 0.5 0.5 • Test of perturbative QCD ... or which way to compute reality is better 50 Q4 flavor scheme 9 0000









→ Study Z+bb production as function of number of b-jet , 1 or 2

Ludivine Ceard, CP3, Petnica 2014

MadGraph

Zbb Signal and backgrounds



Partial Luminosity 2011 : 2.1 fb⁻¹ at 7 TeV Also an analysis with full Lumi 2011 5.0 fb⁻¹ but more complicated.

Signal:

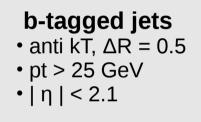
Z + 1 b = one Z + exactly 1 b-jet Z + 2 b = one Z + at least 2 b-jets

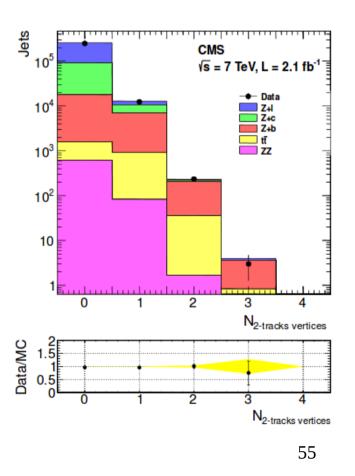
(exclusive) (inclusive) = 1 step in the analysis = next step in the analysis

Phase space:

Z → I^+I^- • I = µ/e, isolated • pt > 20 GeV

• | η | < 2.4





3 Backgrounds:

- Z+c and Z+I : because I and c can be confused with b!
- $t\bar{t} \rightarrow W(l,\nu)b W(l,\nu)b$ same FS than signal + MET
- $ZZ \rightarrow II + bb$



Technically :data



This reconstruction has been done in several steps in T0 then checked by poor shifters, validated, moved to T1 \rightarrow many people are involved in the collection and creation of one dataset.

CMS provides centrally sets of data for which all the events have been triggered by the same characteristic : we know that in any case we want 2 electrons or 2 muons \rightarrow we are going to use the dataset triggered by DoubleEle trigger (with a certain energy so that they are nicely reconstruct), and DoubleMu.

In this dataset = for one event all the objects present have been reconstructed : object electrons and muons with all their characteristics : pT, η , ϕ , E jets reconstructed with different algorithms and all char.

First step = create a skim of this huuuge dataset by putting some of your mandatory requirement in a code : here we want a Z so the combination of the 2 leptons must have an invariant mass around the nominal mass of the Z (90 Gev) \rightarrow 60 and 120 GeV. We can also here ask for one jet. And we can start to identify the jets that can be b jets.

Use the Grid to send your jobs on T2

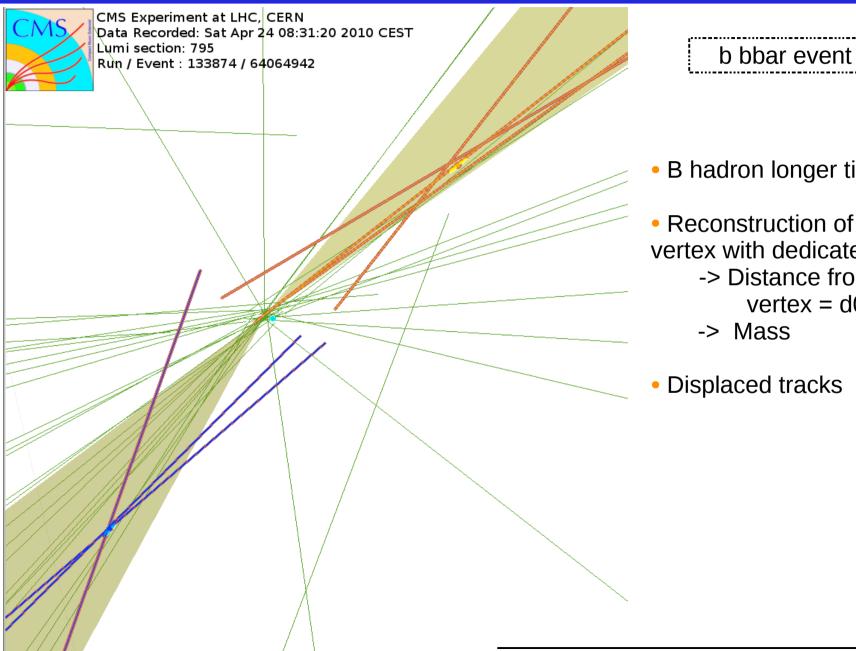
$$(Wc^{2})^{2} = \left(\sum E\right)^{2} - \left\|\sum \mathbf{p}c\right\|^{2}$$
where
$$W \text{ is the invariant mass of the system of particles, equal to the mass of the decay particle.}$$

$$\sum E \text{ is the sum of the energies of the particles}$$

$$\sum \mathbf{p} \text{ is the vector sum of the momentum of the particles (includes both magnitude and direction of the momenta)} 56$$

Zbb : b-tagging





• B hadron longer time of flight

 Reconstruction of a second vertex with dedicated algorithm -> Distance from primary vertex = d0

-> Mass

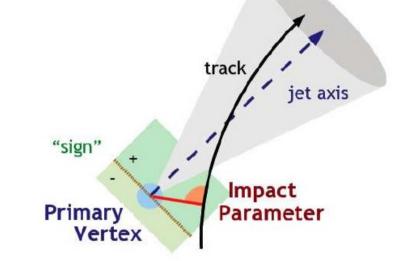
Displaced tracks

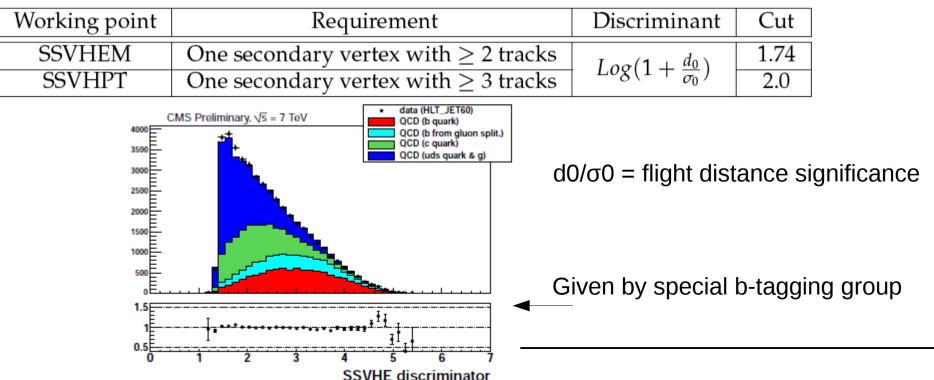




58

- Track-counting (TC): tracks ordered by IP High-Eff (HE) = 2nd track, High-Pur (HP) = 3rd track.
- Simple Secondary Vertex (SSV): at least
 - 2 (High-Eff) or
 - 3 (High-Pur) tracks in vertex fit.







Technically : MC



Groups of people are obliged to run all those MC simulation (running during weeks or months)

CMS provides centrally MC simulation : remember those contain both Generated events and Reconstructed event (just like if they had passed through the detector).

In this dataset = for one event all the objects present have been reconstructed : object electrons and muons with all their characteristics : pT, η , ϕ , E jets reconstructed with different algorithms and all char.

First step = run the exact same code than for data before to create a skim of this huuuuge dataset : you will end up with the same requirement on the Reco objects than for the data.



Use the Grid to send your jobs on T2

Use PhD students to send the jobs and babysit them :)

Remember :you have to do it <u>for the signal and all the backgrounds</u> : We used here Z + jets for : the signal Zb And 2 backgrounds : Z+I and Z + c

We used a ttbar sample for ttbar background.

And a ZZ sample for the ZZ background.

When all your jobs are finished and nice :

Step 2 = putting in place the analysis code with all the cuts and plot all the interesting variables!

- Z → I⁺I⁻
- 2 I = μ/e , isolated
- pt > 20 GeV
- | η | < 2.4
- 60 < mll < 120

- 1 or 2 b-tagged jets
- pt > 25 GeV
- | η | < 2.1
- From algo with HEfficiencies
- No overlap with leptons from the Z : dR > 0.5

10 CEST

Since we run now on the smaller skim : possible to do it on the university local server or Cern server.

Superimpose the data and the all the different MC : signal and backgrounds.

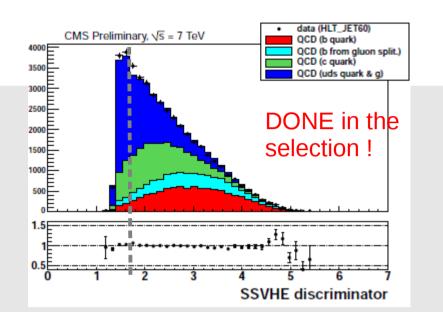
Wait a second : did we get rid of all the background possible ? \ldots mhum \ldots

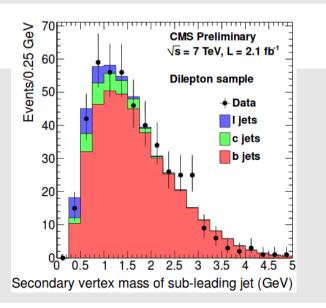
Background: Z + light jets

Z+c, Z+udsg

4

- b-tagging: background reduction
- Detached secondary vertex
- High efficiency selection: 55 %
- 1% mistag





background estimate

- Template fit to the Secondary Vertices Mass
- \rightarrow b purity

We evaluate what we couldn't remove with templates of shapes.

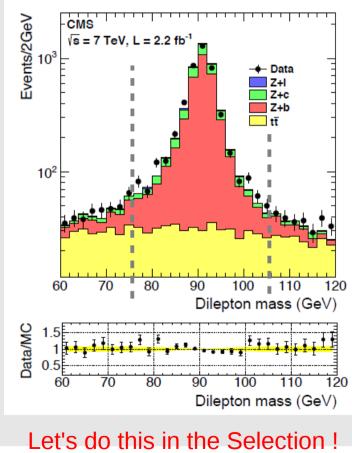


Background: ttbar

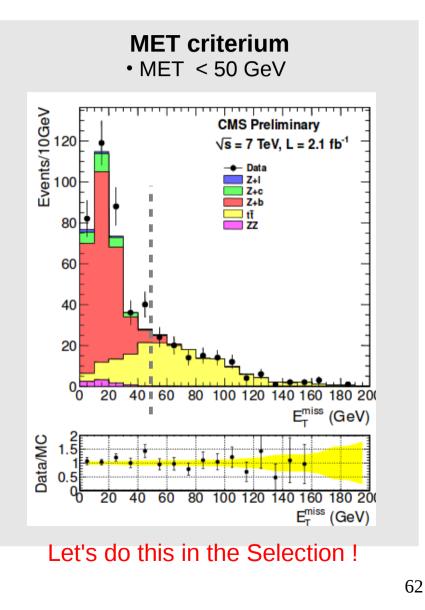


• ttbar reduction:

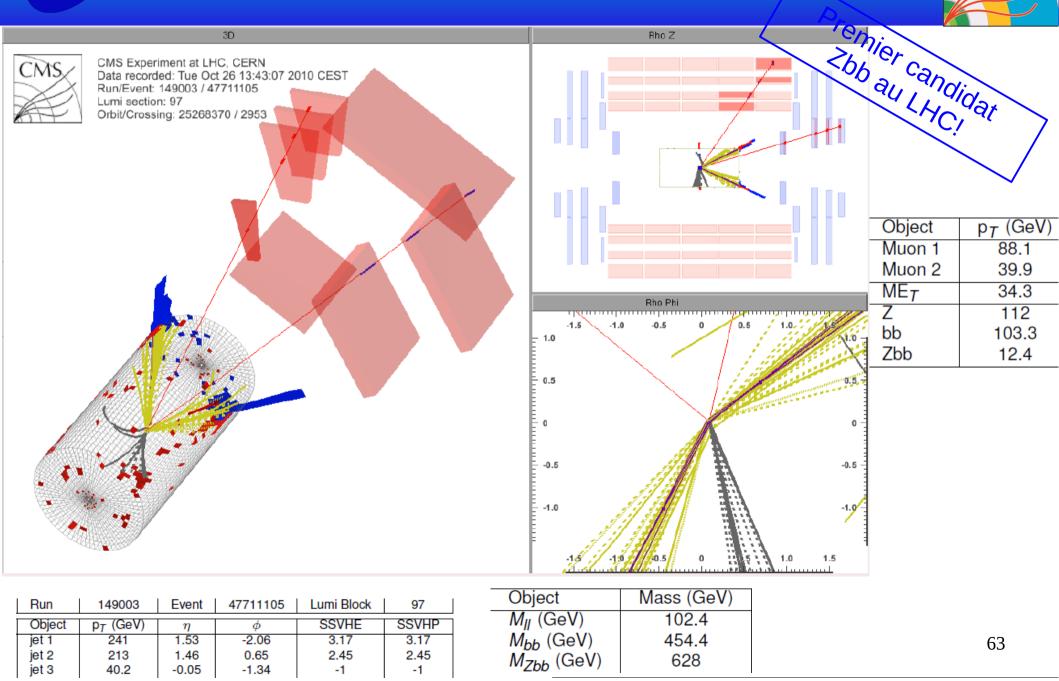
Z mass window • 76 < m_{_} < 106 (GeV) CMS √s = 7 TeV, L



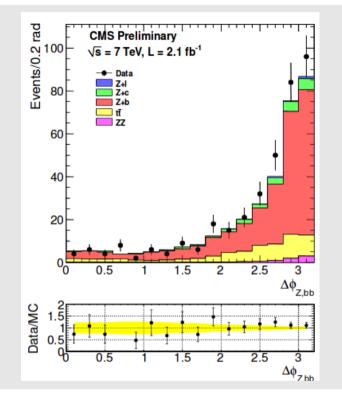
 \bullet ttbar estimate: fit to $m_{_{I\!I}}$ from templates



And finally : an event!



And finally : The plots!

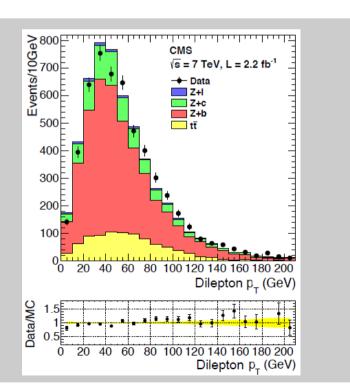


Momentum distributions :

- Slightly harder spectrum for data than MC at LO
 - → NLO to be considered

Angular distributions :

- General reasonable agreement data/MC (MadGraph)
- MPI contribution reasonably modeled from low values of $\Delta\phi_{Z,bb}$

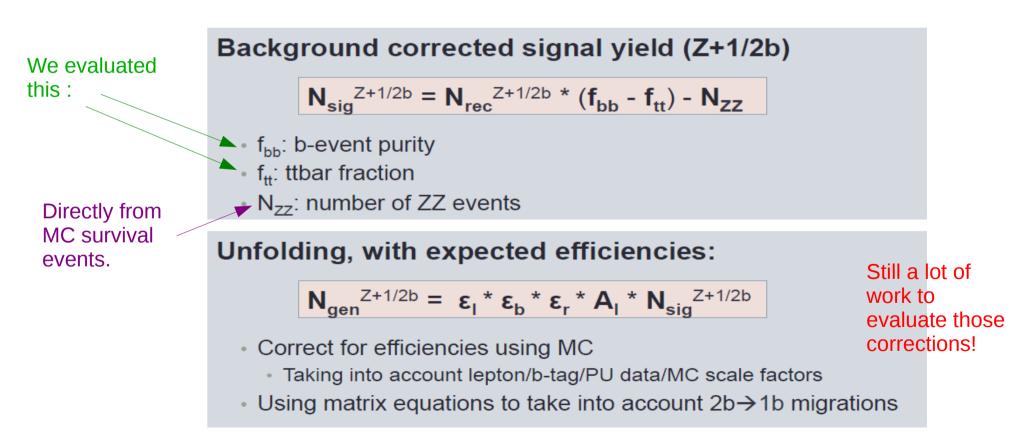


Ludivine Ceard, CP3, Petnica 2014

4

Calculation of a σ

We need to have access at the number of Signal events = Number of data selected – number of background events we couldn't eliminate



Unfolding and Results



Unfolding:

reconstructed b-jets → # hadron-level b-jets

$$\begin{pmatrix} \sigma(Z+1b) \\ \sigma(Z+2b) \end{pmatrix} = \frac{1}{\mathcal{L}} \times \mathbf{A}^{-1} \times \mathbf{E}_{r}^{-1} \times \mathbf{E}_{l}^{-1} \times \mathbf{E}_{b}^{-1} \times \begin{pmatrix} N_{sig}^{Z+1b} \\ N_{sig}^{Z+2b} \end{pmatrix}$$

Corrections for all efficiencies and acceptance

Cross sections:

• Results for ee and $\mu\mu$ channels compatible and combined in a single measurement

Cross sections at the particle level		
Multiplicity bin	Combination	
$\sigma_{hadron}(Z+1b, Z \rightarrow II)(pb)$ $\sigma_{hadron}(Z+2b, Z \rightarrow II)(pb)$	$3.41 \pm 0.05 \pm 0.27 \pm 0.09$ $0.37 \pm 0.02 \pm 0.07 \pm 0.02$	
$\sigma_{hadron}(Z+b, Z \rightarrow II)(pb)$	$3.78 \pm 0.05 \pm 0.31 \pm 0.11$	

Measurement: $\sigma(Z(II)+2b) = 0.37 \pm 0.02 \text{ (stat.)} \pm 0.07 \text{ (syst.)} \pm 0.02 \text{ (theory) pb}$ MadGraph expectation: $\sigma(Z(II)+2b) = 0.33 \pm 0.01 \text{ (stat) pb}$

• Compatible with expectations from MadGraph 5 flavor corrected to NNLO

66

Systematic errors on the measurement

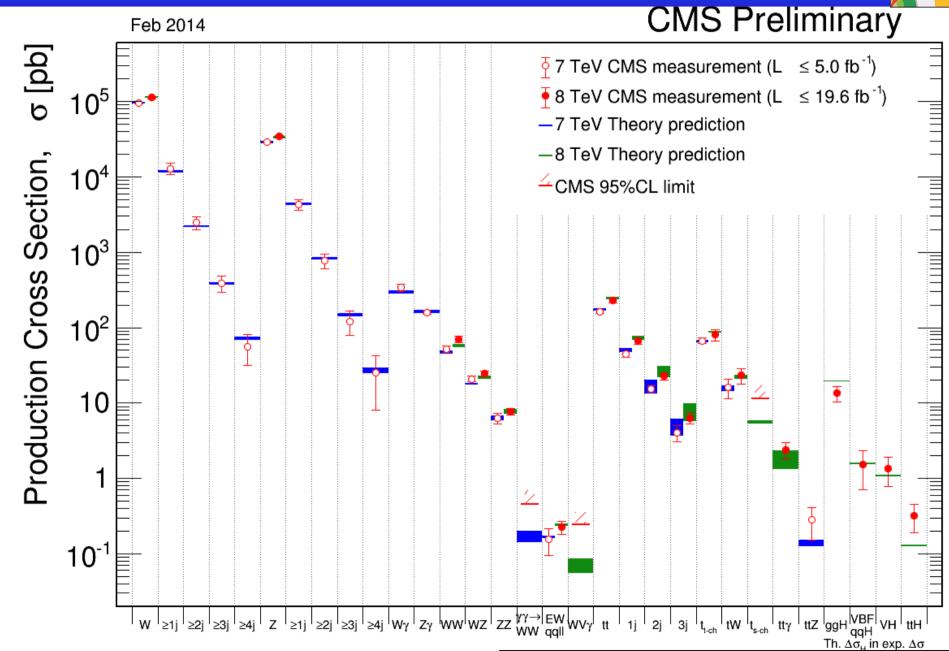


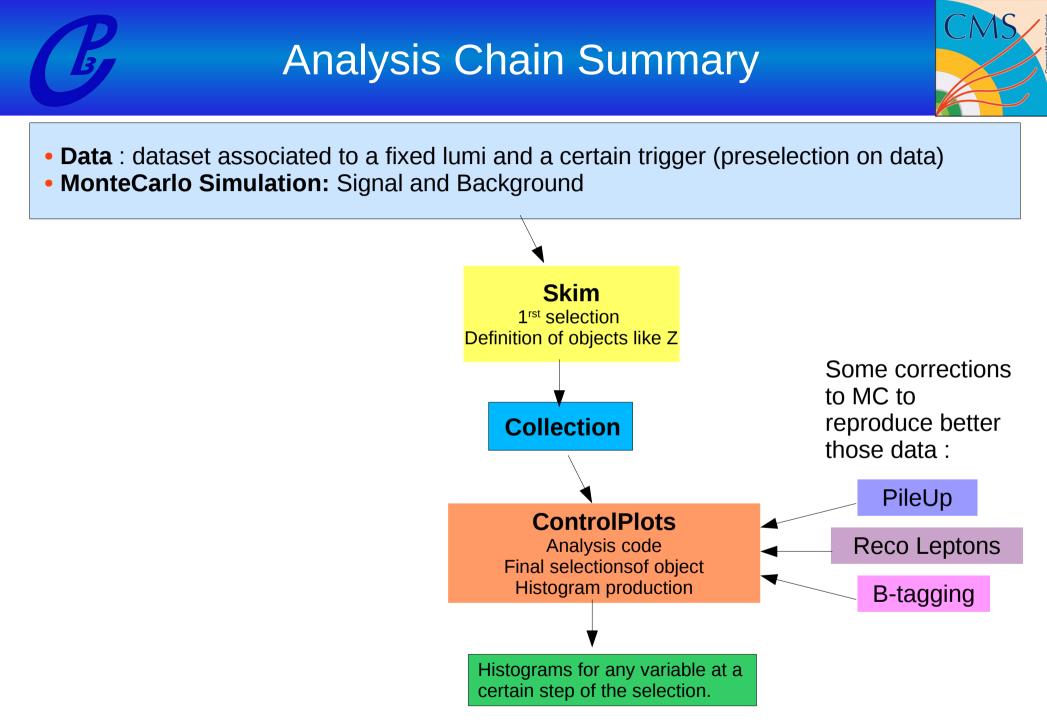
	Correlated sources	Fractional uncertainty (%)		
	b-tagging efficiency	10		
	b-jet purity	5.6 (ee+b)	4.6 (μμ+b)	
	tt contribution	2.9		
	Jet energy scale	2.5		
	Luminosity	2.2		
	Jet energy resolution	0.5		
0.6 interactions	Pile-up	1.5 (ee+b)	0.5 (μμ+b)	
	Mistagging rate	0.04		
	Theory (via \mathcal{A}_{ℓ})	$+4.2 \\ -6.5$		
	Theory (via \mathcal{C}_{hadron})	$^{+0.7}_{-6.9}$		
	Uncorrelated sources	ee+b	μμ+b	
	Trigger and dilepton selection	4	2	
	tt contribution	1.9	2.2	
	Experimental systematic	13.0	12.3	
	Theoretical systematic	$+4.2 \\ -9.5$	$^{+4.2}_{-9.5}$	
	Statistical	2.2	1.7	

±

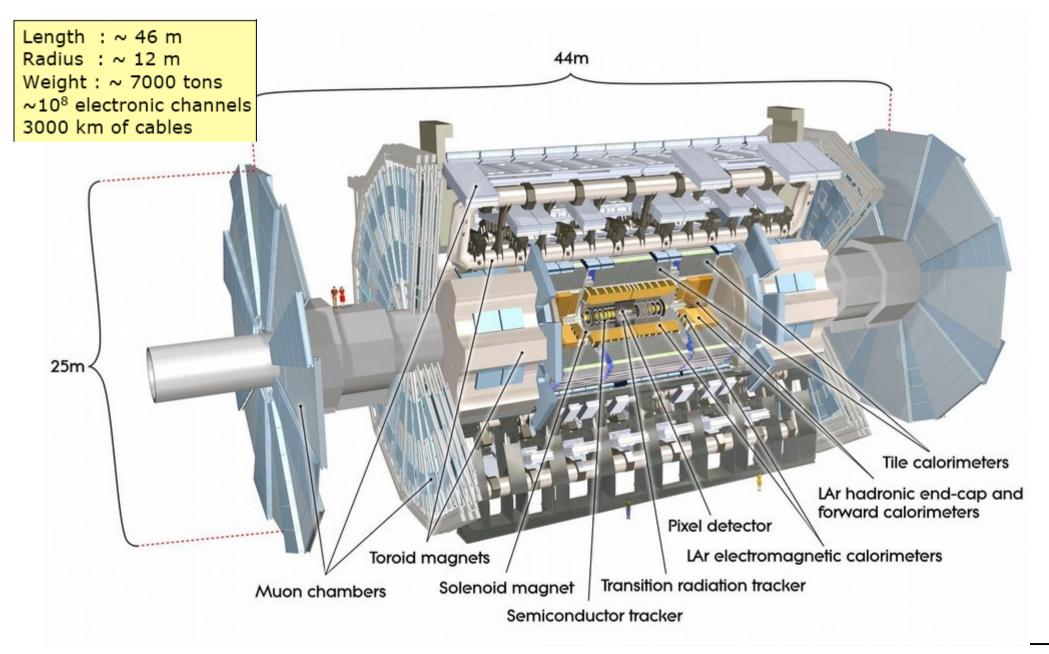








ATLAS ATLAS : A Toroidal LHC ApparatuS





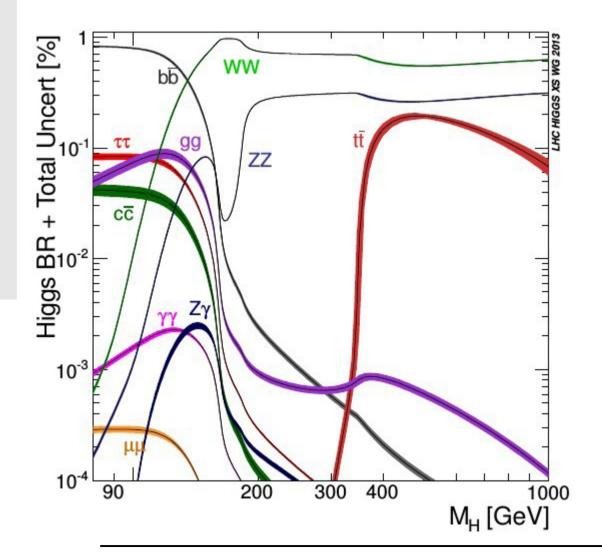
ATLAS + Higgs Search (and found!)

 $\textbf{H} \ \rightarrow \ \textbf{ZZ} \ \rightarrow \ \textbf{4leptons}$

One of the most sensitive channels.

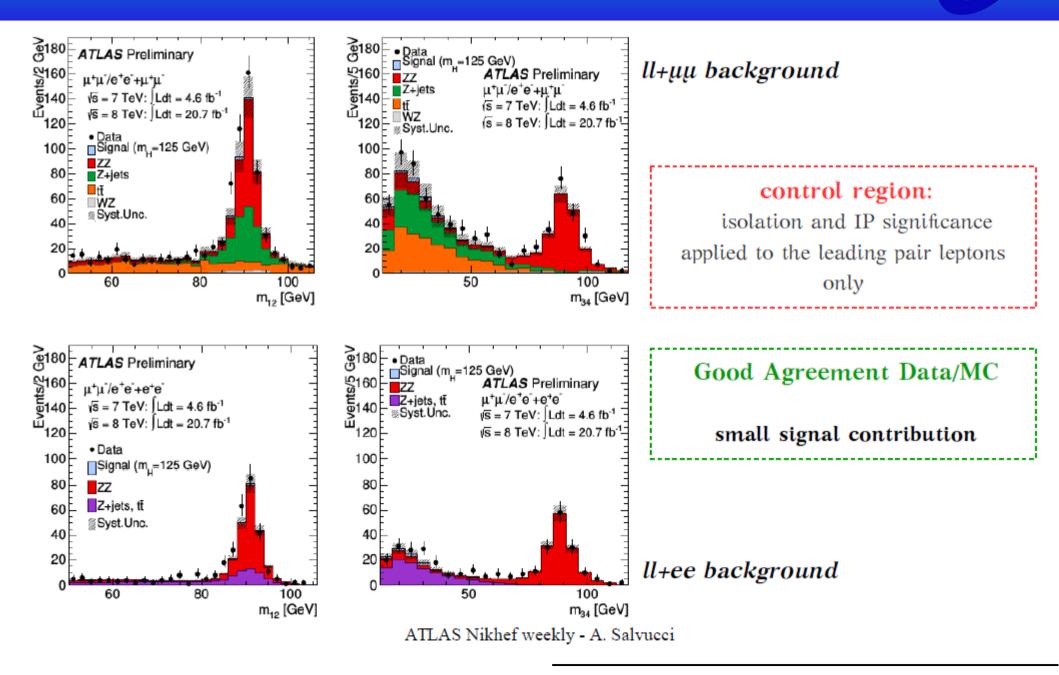
It provides a rather clean final state signature.

Final state fully reconstructed Best mass resolution Low BR fraction (at low mass)





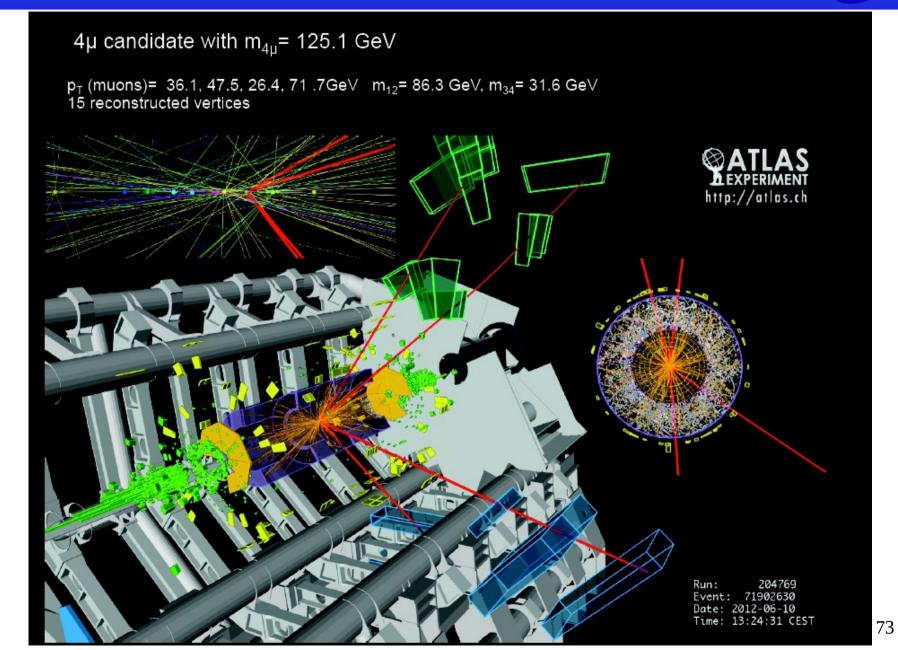
H-> ZZ-> 4 leptons





$H \rightarrow 4$ leptons







H-> ZZ-> 4 leptons

Animated gif which is making my computer die!

Announcement the 4th of July 2012

EXPERIME

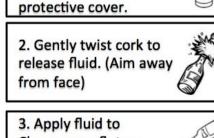




For use in case

o Higgs disco

COURANCE

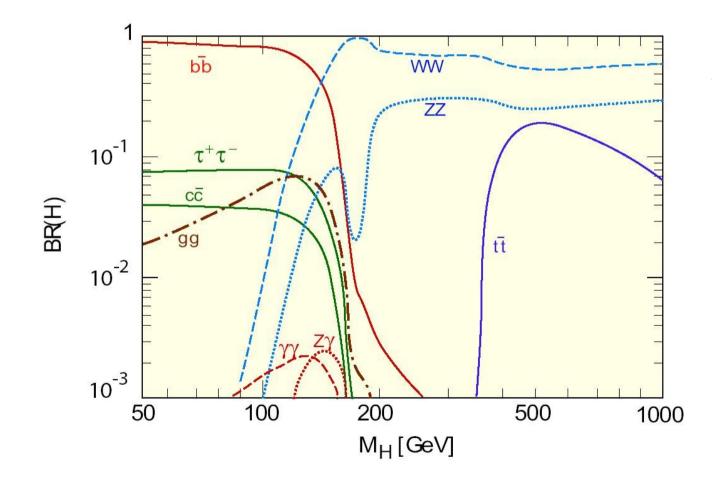


3. Apply fluid to Champagne flutes. Repeat until all flutes are filled.

More on Higgs discovery

How to observe the boson

Observe the Higgs == observe a clear excess with respect to the expected background (predicted for a case with no Higgs)

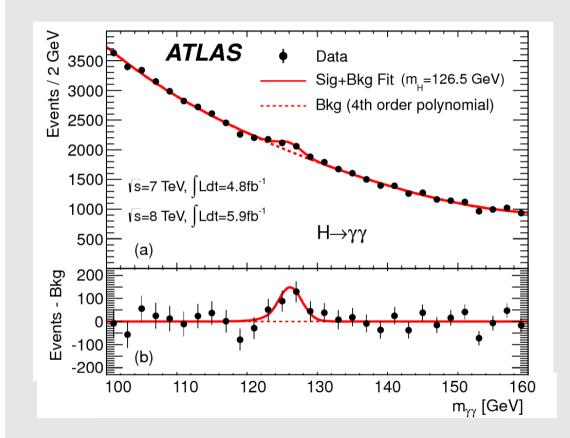


Several decay modes with different probabilities to happen in function of the mass it could have.

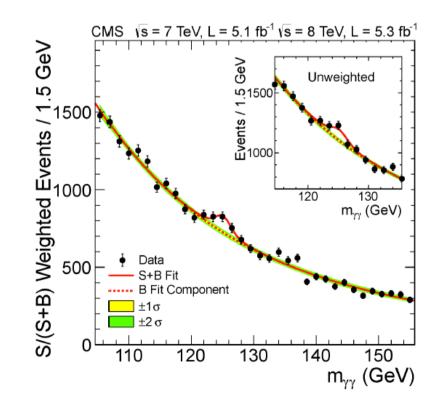




Н-> уу



Deviation from background B. B = all standard model processes that can lead to a yy final state

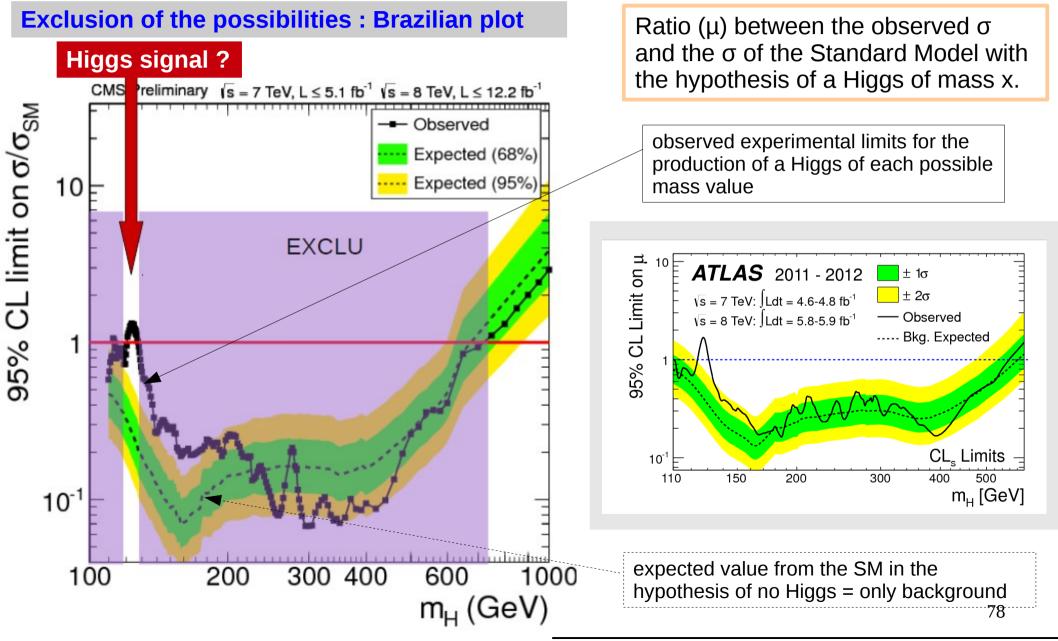


Ludivine Ceard , CP3 , Petnica 2014

77







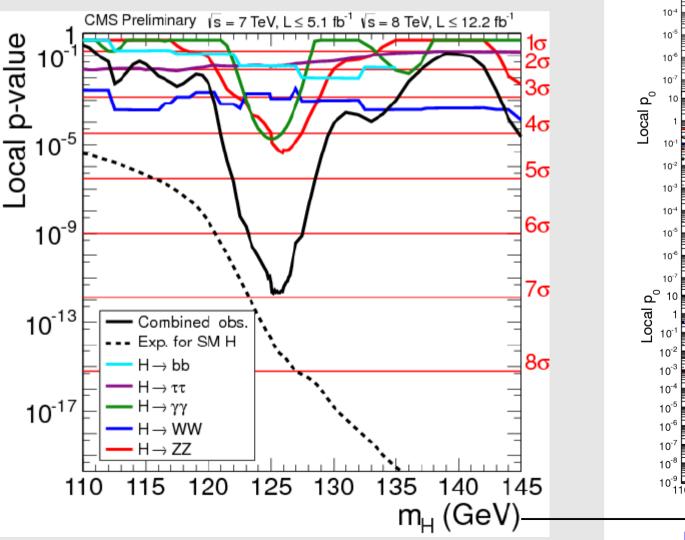
ATLAS EXPERIMENT

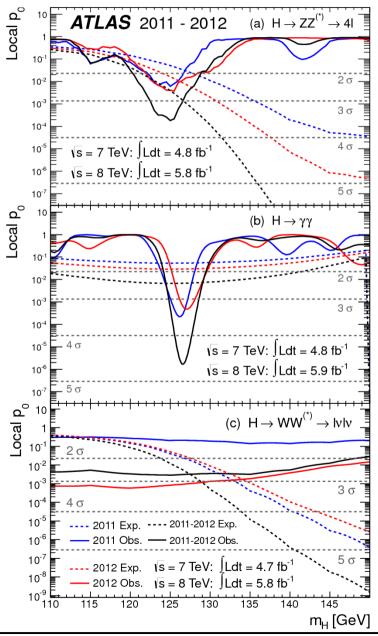
Higgs discovery



Probability : p-value plot



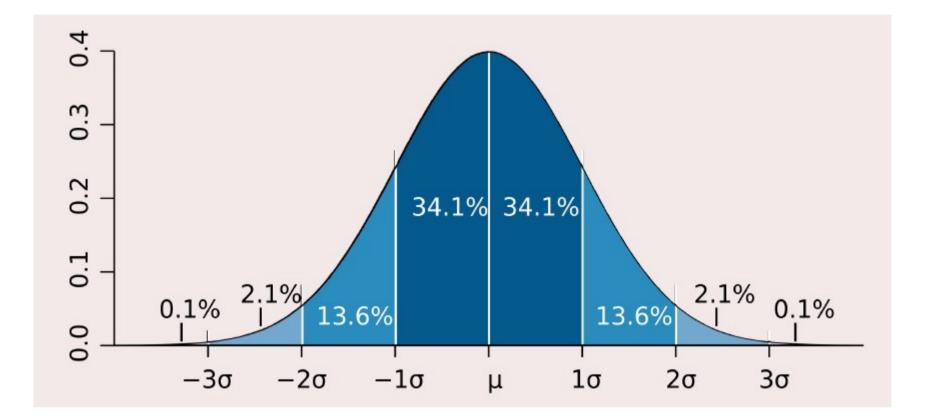






Standard deviation

The number of standard deviation is a convention based on the Gaussian to express the small probabilities



By convention/tradition we talk about evidence for 3σ et of discovery for 5σ .





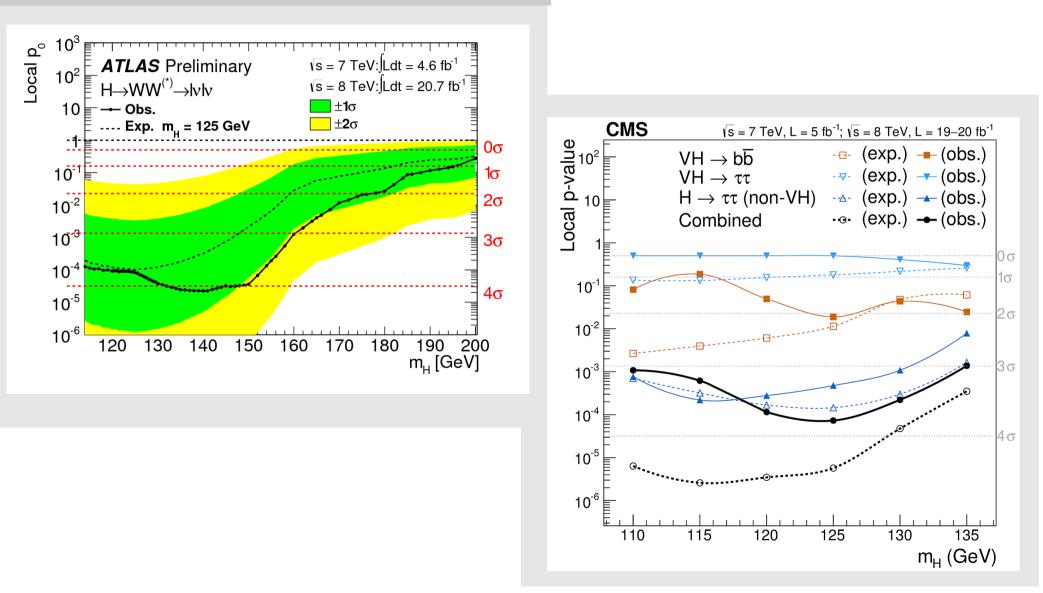
Per channel

Discovery! **Evidences** CMS Preliminary (s = 7 TeV, $L \le 5.1$ fb⁻¹ (s = 8 TeV, $L \le 12.2$ fb⁻¹ CMS Preliminary (s = 7 TeV, $L \le 5.1$ fb⁻¹ (s = 8 TeV, $L \le 12.2$ fb⁻¹ Local p-value Local p-value - -1σ 1σ 2σ 2σ 0-2 10⁻² 3σ 3σ 4σ 4σ 10⁻⁵ I 10⁻⁵ 5σ 5σ 10⁻⁸ 10⁻⁸ 6σ 6σ Combined obs. 10⁻¹¹ 10⁻¹¹ 7σ 7σ --- Exp. for SM H Combined obs. $-H \rightarrow WW$ Exp. for SM H 10^{-14} $H \rightarrow bb$ **10**⁻¹⁴ H→γγ $H \rightarrow \gamma \gamma + H \rightarrow Z$ $H \rightarrow bb + \tau\tau +$ 8σ 8σ Η→ττ 130 115 120 125 135 140 145 110 115 120 125 130 135 40 145 1 m_H (GeV) m_H (GeV)

81

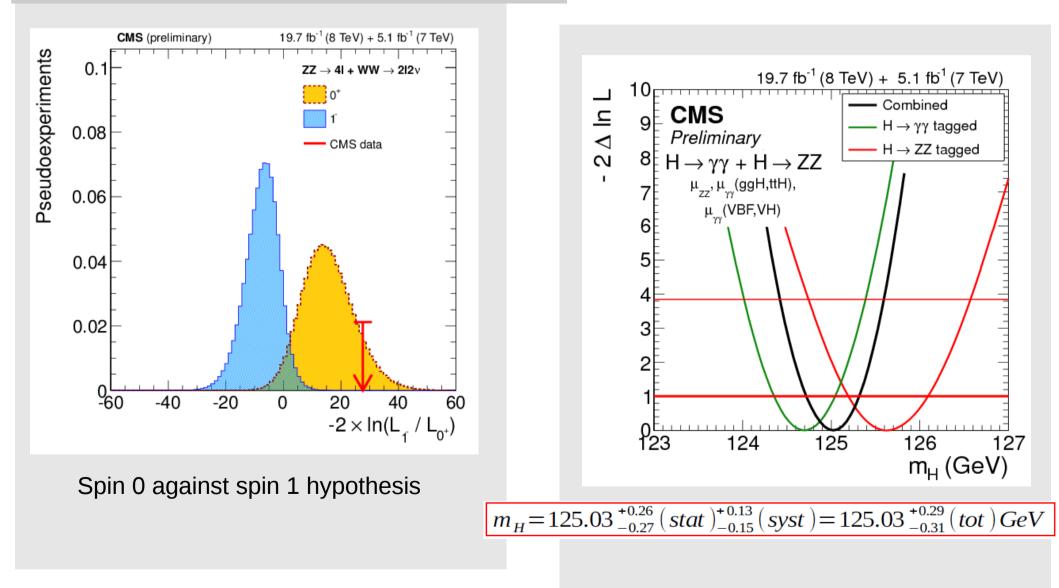


Latest H → WW and to fermions



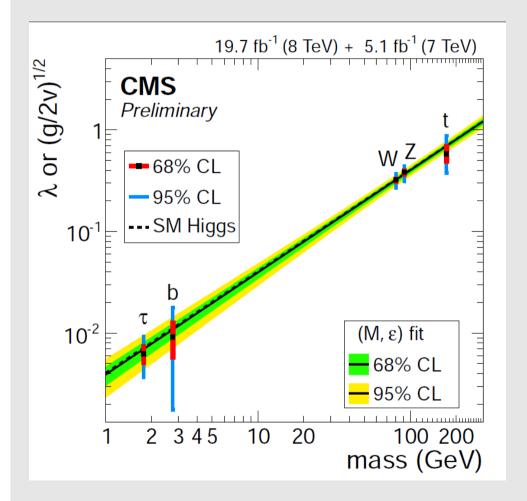


Property measurements :

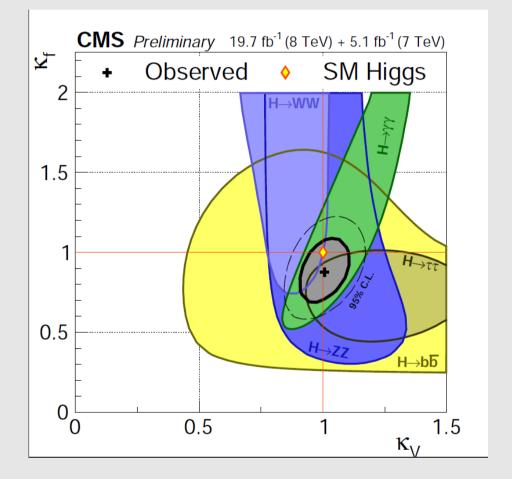




Property measurements :



Coupling strength of Higgs as function of particle mass.



Combination for coupling to fermions and bosons

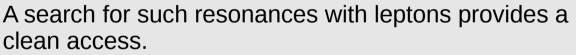
CMS + exotic search (and not found!)

 $\gamma^*/\mathbf{Z}/\mathbf{Z}'$



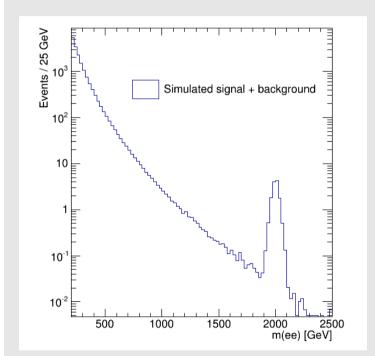
$Z' \rightarrow 2leptons$

- Many theories predict heavy resonances.
 - Extra dimensions heavy Z's (spin 1) heavy gravitons (spin 2)
 - Grand unified theories



Best performance with electrons and muons. Decaying taus are more difficult to handle.

Leptons have high momentum and care is necessary for the selection



Search for a narrow resonance (Detector resolution much wider than natural width of the resonance) on steeply falling SM background

86

CMS + exotic search (and not found!)

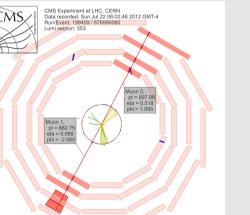
Z' → 2leptons

Electron selection:

Two electrons candidates with high energy deposit in the ECAL and associated track.

No other energy deposits around the electron candidate (ECAL isolation).

No other tracks in a cone around the electron candidate (Tracker isolation).

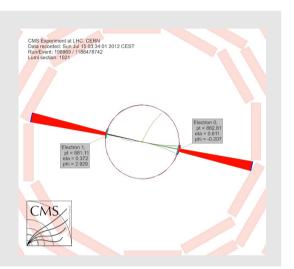


Muon selection:

Two opposite charge muon candidate tracks reconstructed in the inner tracker and the muon chambers.

No other tracks in a cone around the muon candidate (Tracker isolation).

Both tracks must be close to the same vertex (Cosmic muon rejection).



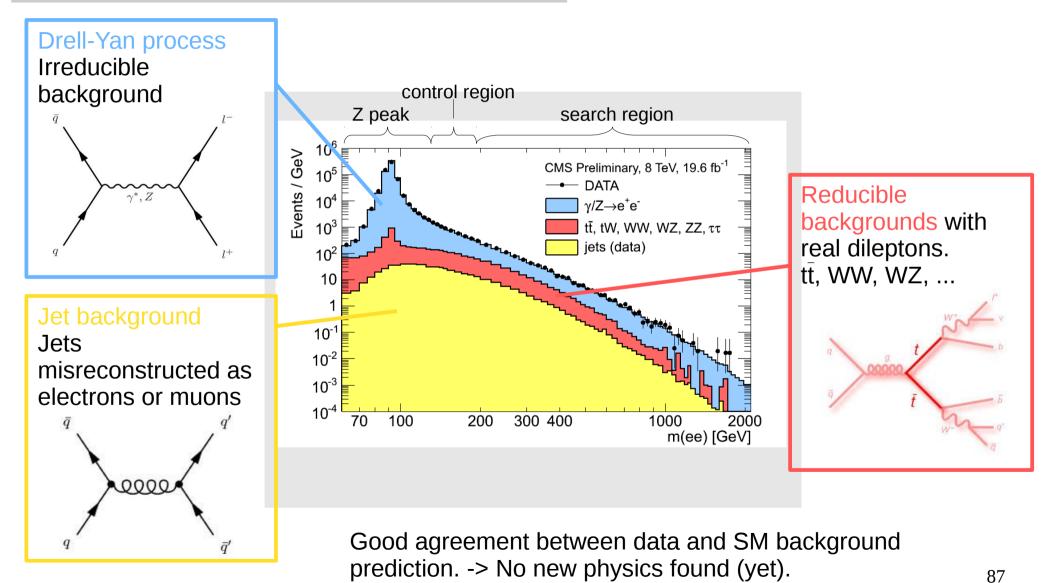




Z' → ||



Result Invariant Mass :



Z' → ||



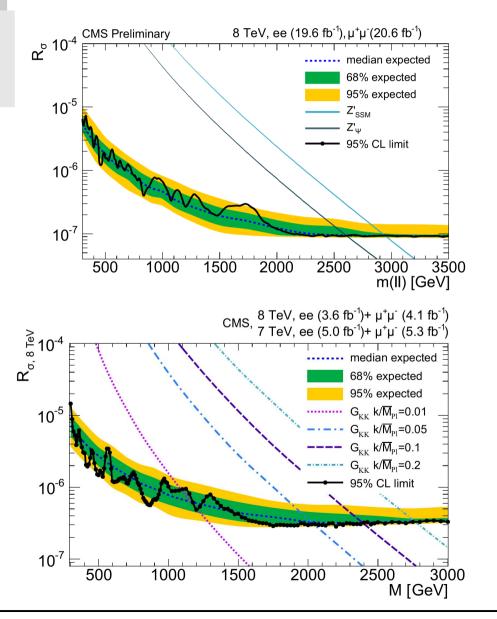
In absence of a significant excess over the SM we set limits on the cross section of new physics processes.

• Spin 1 and spin 2 resonances have different kinematics -> different acceptance in the detector -> Different limit plots necessary

• Limit curve below the theory line -> theory excluded at > 95% confidence level

• Intersection between limit curve and theory curve give the mass up to which the theory is excluded at 95% CL

In the spin 1 case: Z'_{SSM} excluded up to 2.96 TeV Z'_{ψ} excluded up to 2.6 TeV

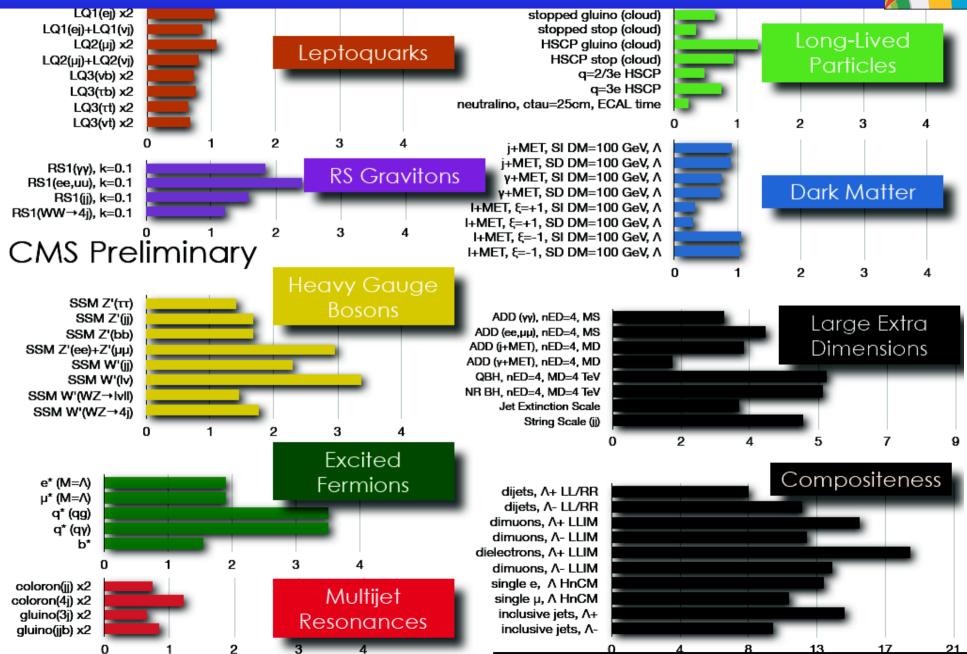






Exotica status

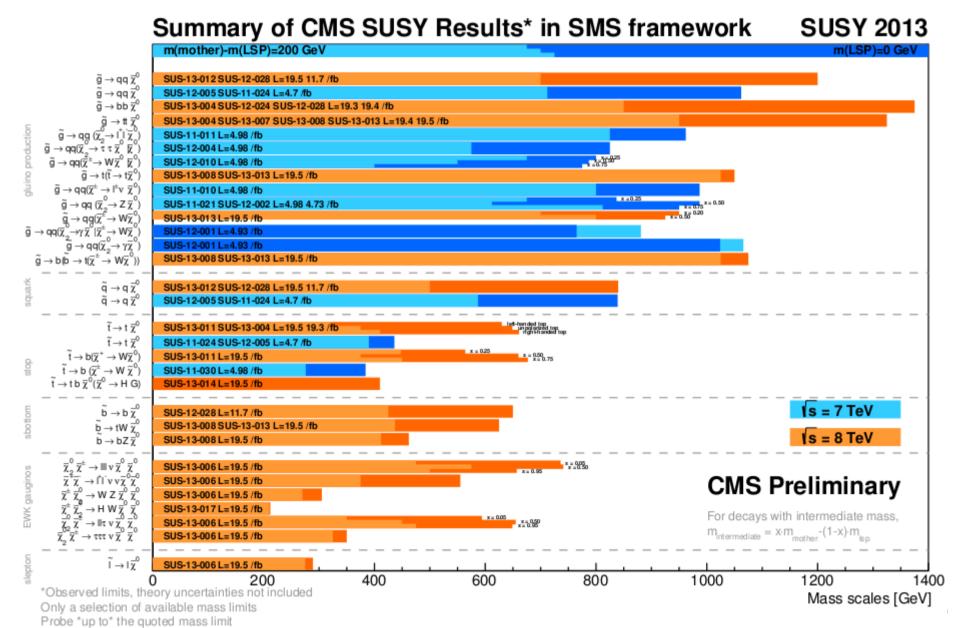


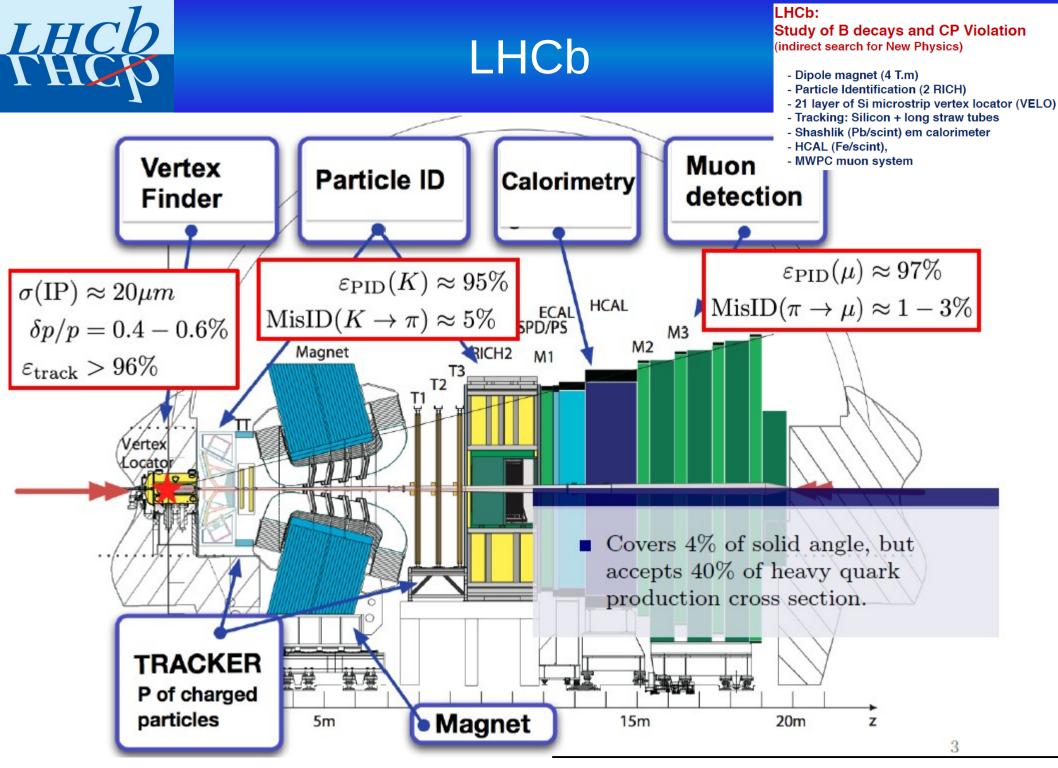




SUSY state



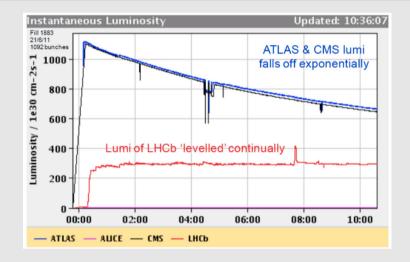








Luminosity levelling

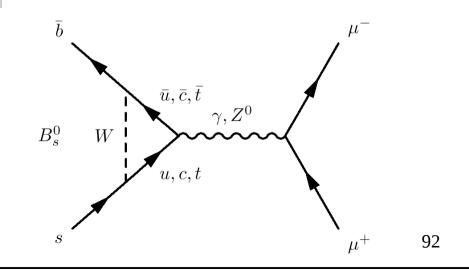


- LHCb designed to run at lower luminosity than ATLAS/CMS.
- LHCb tracking/PID is sensitive to pile-up.
- → LHC pp beams are displaced to reduce instantaneous luminosity stable running conditions.
- L 2011 : 2.7* 10³²cm⁻²s⁻¹
- L 2012 : 4.0* 10³²cm⁻²s⁻¹

Searching for new physics

INDIRECT search : if it's not a W but a heavier (unknown) particle in the virtual quantum loop Branching Ratio $B_s \rightarrow \mu\mu$ will be different than the one expected from the standard model :

$$\mathcal{B}(B_s^0 \to \mu\mu) = (3.35 \pm 0.28) \times 10^{-9} \mathcal{B}(B^0 \to \mu\mu) = (1.07 \pm 0.10) \times 10^{-10}$$



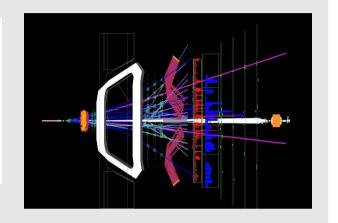


Bs μμ

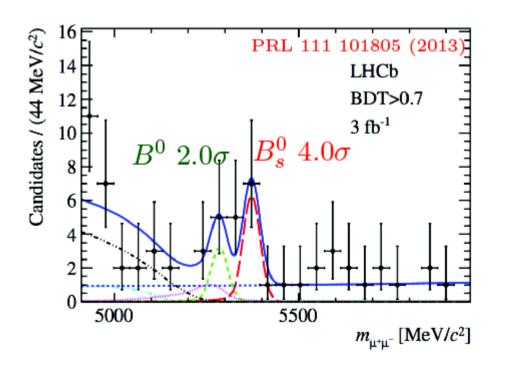


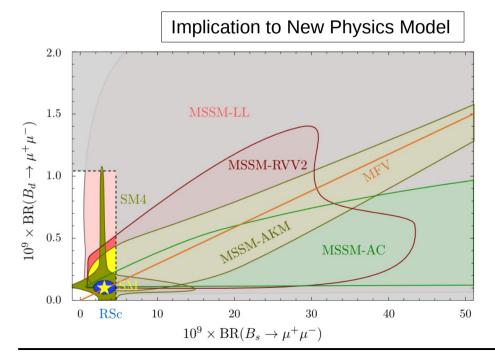
First evidence of $B_s \rightarrow \mu\mu$:

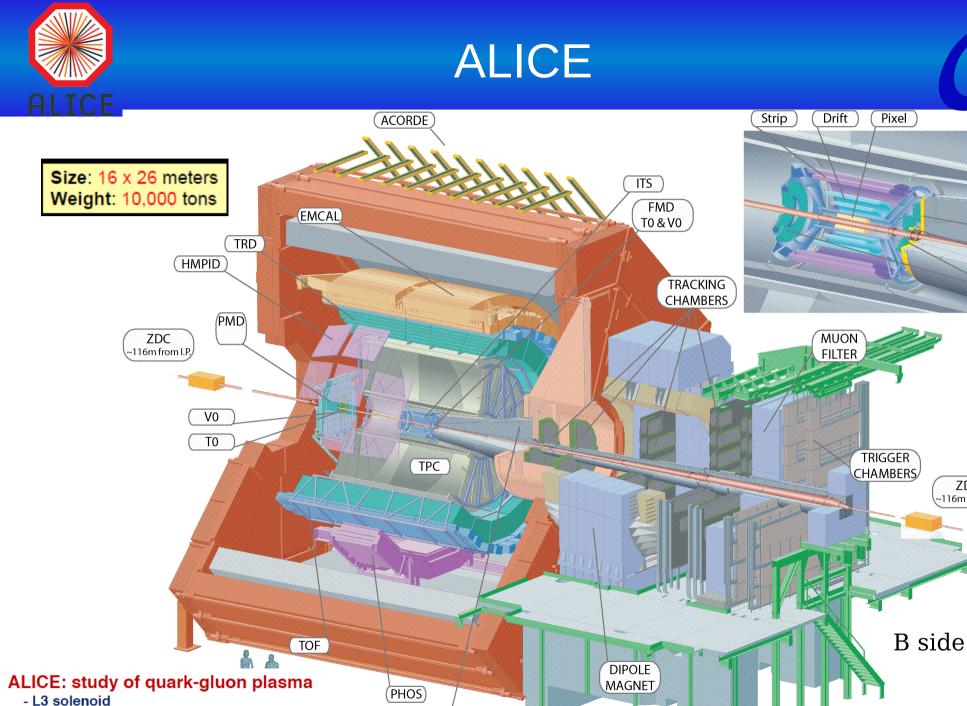
	$\mathcal{B}(B_s^0 \to \mu\mu) \times 10^{-9}$	$\mathcal{B}(B^0 \to \mu\mu) \times 10^{-9}$
LHCb	$2.9^{\pm1.1\pm0.3}_{\pm1.0\pm0.1}$	$3.7^{+2.4+0.6}_{-2.1-0.4}$
\mathbf{CMS}	$3.0^{+1.0}_{-0.9}$	$3.5^{+2.1}_{-1.8}$
Combined	2.9 ± 0.7	$3.6^{+1.6}_{-1.4}$



Compatible with the SM predictions.







ABSORBER

- Large TPC
- Si microstrip, drift and pixels detectors - Particle identification: RICH, TRD, TOF
- PbWO₄ crystals + Pb/scintillator ecal
- Single arm forward muon system

Ludivine Ceard, CP3, Petnica 2014

V0

Τ0

FMD

ZDC -116m from I.P.

с .



ALICE Physics



Pb+Pb @ sqrt(s) = 2.76 ATeV

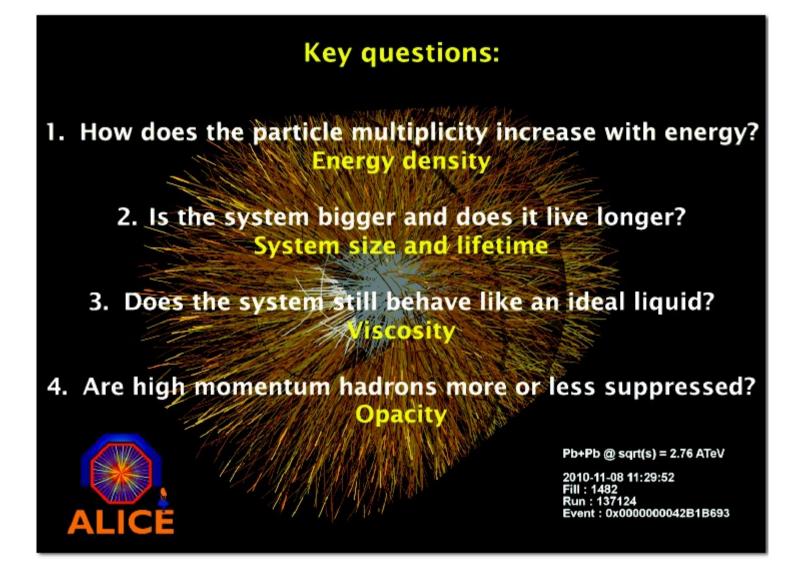
2010-11-08 11:30:46 Fill : 1482 Run : 137124 Event : 0x00000000D3BBE693

95



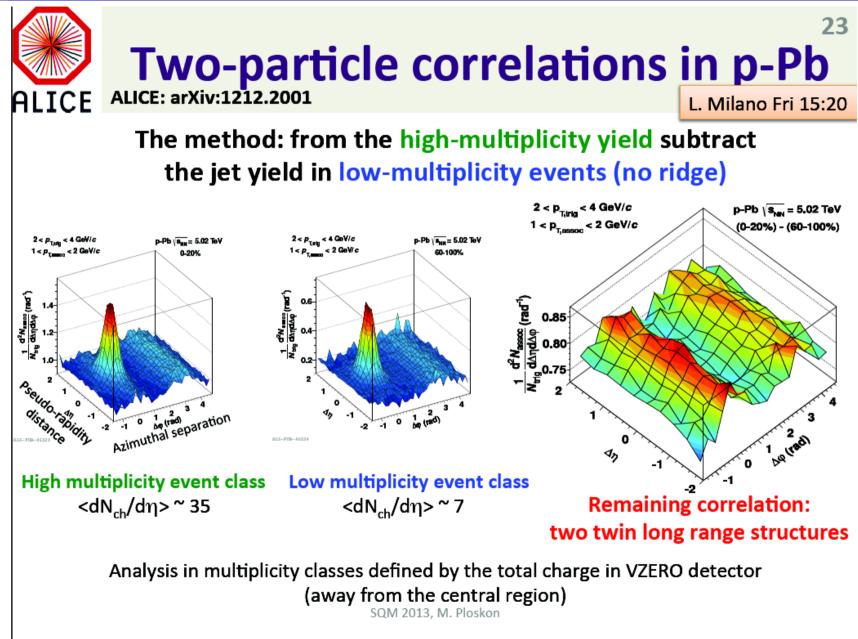
ALICE Physics







ALICE Physics



97

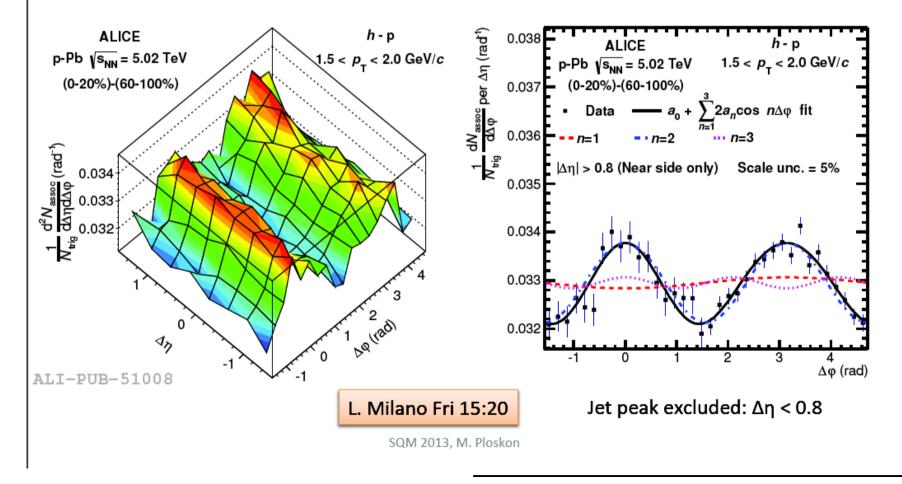


ALICE

ALICE Physics

Twin ridge structure in p-Pb with identified particles

Shown here: hadron-proton correlation (high-low mult. percentile subtracted)



98

25

Other experiments at CERN

Many !!!!!!

Using the LHC :

- Totem forward to CMS
- LHCf forward to Atlas
- AD : from PS beam :Antiproton Decelerator : studying antimatter, first anti H atom!
- Gran Sasso (Italy) : using the neutrinos from SPS (and sometimes see them faster

than speed of light)

What is TOTEM?

TOTEM will measure the effective size or 'cross-section' of the proton at LHC. To do this TOTEM must be able to detect particles produced very close to the LHC beams. It will include detectors housed in specially designed vacuum chambers called 'Roman pots', which are connected to the beam pipes in the LHC. Eight Roman pots will be placed in pairs at four locations near the collision point of the CMS experiment. TOTEM has more than 70 members from 10 institutes in 7 countries (May 2007).

Size	440 m long, 5 m high and 5 m wide	
Weight	20 tonnes	
Design	roman pot and GEM detectors and cathode strip chambers	
Material cost	6.5 MCHF	
Location	Cessy, France (near CMS)	

What is LHCf?

LHCf is a small experiment that will measure particles produced very close to the direction of the beams in the proton-proton collisions at the LHC. The motivation is to test models used to estimate the primary energy of the ultra high-energy cosmic rays. It will have detectors 140 m from the ATLAS collision point. The collaboration has 21 members from 10 institutes in 6 countries (May 2007).

Size	two detectors, each measures 30 cm long,	
	10 cm high, 10 cm wide	
Weight	40 kg each	
Location	Meyrin, Switzerland (near ATLAS).	

Give a break to LHC :

- ISOLDE
- Cloud, NA48, CAST, COMPASS

For more information visit: http://totem.web.cern.ch/Totem/





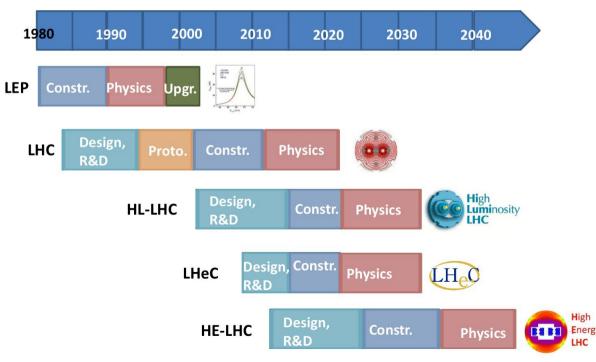
Conclusion



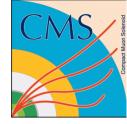
• LHC has fulfilled his contract of discovery machine :

A new scalar particle very SM-Higgs-boson-like has been discovered!

- But for the moment nothing that haven't been predicted by SM : Where is the New Physics?
- All physicists are waiting for the restart of the LHC in 2015 to keep on looking!





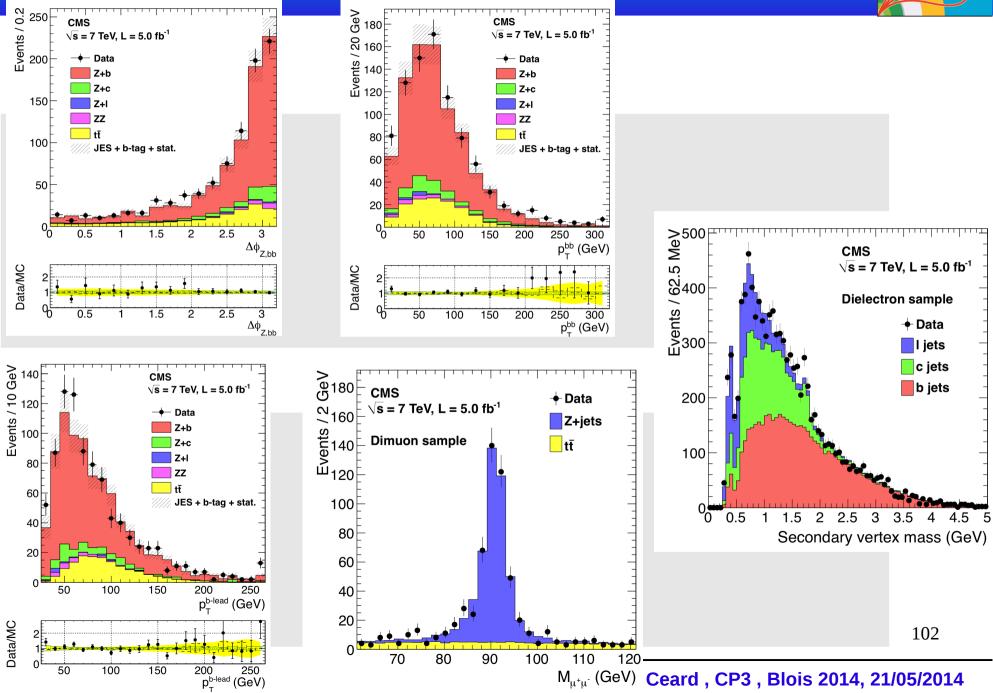


Back-up

ATLAS EXPERIMENT

Z + b











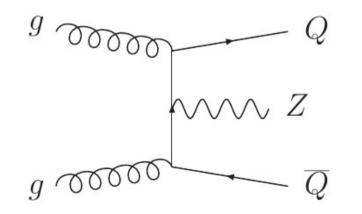
Calculations currently derived in 2 ways:

4-flavour scheme

Dittmaier, Kramer, Spira, Dawson, Jackson, Reina, Wackeroth

Explicit gluon splitting -> divergences if massless b.

massive b, g -> bb , 1 or 2 b observed

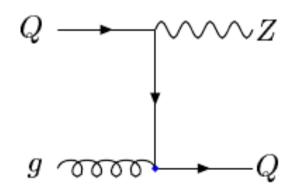


5-flavour scheme

Campbell, Ellis, Maltoni, Willenbrock

g -> bb inside b-PDF -> all orders = no divergences.

in calculations, b is massless. second b added during parton shower and hadronisation by Pythia

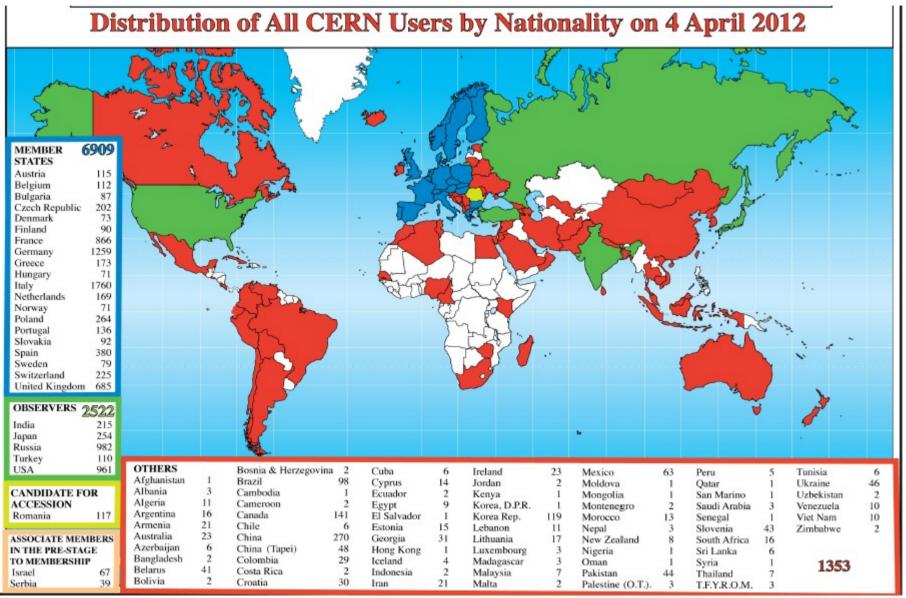


Should agree at NLO

LHC Run



Global Science

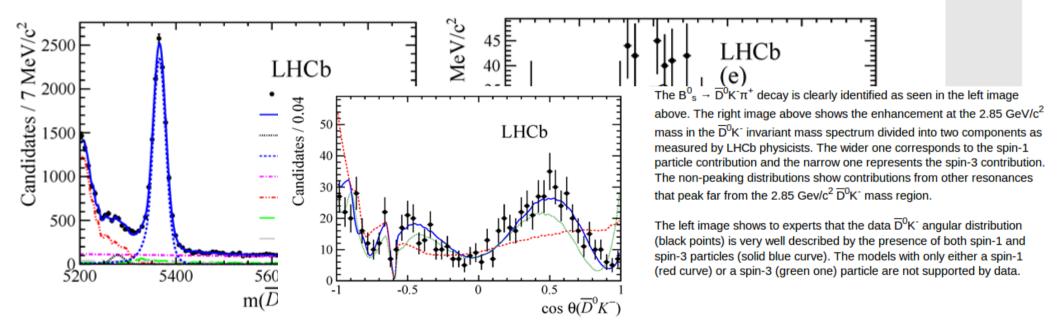




LHCb Other results

15 July 2014: First observation of a heavy flavored spin-3 particle

Today at the 15th International Conference on B-Physics at Frontier Machines at the University of Edinburgh, <u>Beauty 2014</u>, the LHCb collaboration has presented the results of a study of strange beauty meson B^0_s decay into an anti-charm meson \overline{D}^0 , a K⁻ meson and a π^+ meson ($B^0_s \rightarrow \overline{D}^0 K^- \pi^+$). Previous results indicated the existence of a strange-charm $D^*_{sJ}(2860)^-$ particle in the $\overline{D}^0 K^-$ invariant mass spectrum, and the study of the $B^0_s \rightarrow \overline{D}^0 K^- \pi^+$ decay allows one to study this structure and measure its properties. Today's LHCb observation shows with 10 σ significance that, in fact, this excess seen in the $\overline{D}^0 K^-$ mass spectrum is composed of two particles with different spins, spin-1 and spin-3. This is the first observation of a heavy flavored spin-3 particle, and the first time that any spin-3 particle has been seen to be produced in B decays.



click the images for higher resolution

The $D_{sJ}^*(2860)^-$ particles are composed of an anti-charm quark \overline{c} and a strange quark s. The quark-anti-quark pair is bound by strong interactions and can form different quantum states with different values of spin and angular momentum in analogy to the different quantum states of ordinary atoms. The presence of the spin-3 contribution gives a clear signature that both particles are members of the so called 1D family having two units of angular momentum between the quark and the antiquark. This discovery demonstrates that the spectroscopy of the 1D families of heavy flavoured mesons can be studied experimentally. Further insights can be expected with similar analysis of B decays at LHCb and the LHCb upgrade.

LUUIVIILE CEALU, CFS, FELIICA 2014



LHCb Other results

14 November 2011: CP violation in charm decays.

$[\Delta A_{CP} = (-0.82 \pm 0.21 \pm 0.11)\%]$

The LHCb Collaboration has presented today at the <u>Hadron Collider Particle Symposium</u> in Paris possible first evidence for CP violation, the difference between behaviour of matter (particles) and antimatter (antiparticles), in charm decays. The study of CP violation in both charm and beauty particle decays is central to the LHCb physics programme. In the Standard Model CP violation is expected to be very small in the charm sector, whereas new physics effects could generate enhancements.

In this new analysis the LHCb physicists have used data collected in the first half of the 2011 run to study the differences in decay rates of neutral D meson particles composed of a <u>charm</u> quark c bound with an <u>up</u> antiquark (\vec{u}) and D meson antiparticles (\vec{D}) composed of a <u>charm</u> antiquark (\vec{c}) bound with an <u>up</u> quark (u). The decays of D^{*+} mesons into D mesons and <u>m</u>⁺, and D^{*-} mesons into D mesons and <u>m</u>⁻ were used to select the D and D mesons. In the next step of the analysis the difference (asymmetry A_{CP}) between the decay rates of D and D mesons into K⁺K⁻ pairs as

well as into $\pi^+\pi^-$ pairs was measured. By determining the *difference*, ΔA_{CP} , in CP asymmetries for the K+K- and $\pi^+\pi^-$ decays, the analysis strongly suppresses possible measurement biases which could arise through effects related to particle production, selection etc. The following preliminary result is obtained:

 ΔA_{CP} = (-0.82 ± 0.21 (stat.) ± 0.11 (sys.))% [3.5 sigma significance for experts]

A very interesting period now [$x'^2 = (5.5 \pm 4.9) \times 10^{-5}$; y' = $(4.8 \pm 1.0) \times 10^{-3}$] theoretical work will be require [$A_{\Gamma}(KK) = (-0.35 \pm 0.62 \pm 0.12) \times 10^{-3} A_{\Gamma}(\pi\pi) = (0.33 \pm 1.06 \pm 0.14) \times 10^{-3}$]

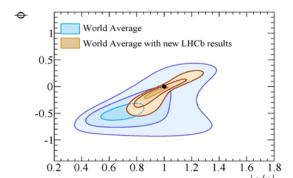
40000 -10000 -10000 -10000 -1820 1840

The LHCb Collaboration has reported recently new important results on charm physics.

(1) Ten months ago, the LHCb Collaboration presented the first observation of the $D^0 - \overline{D}^0$ oscillations in which the D^0 matter mesons turn into their antimatter partners. Contrary to the $B^0 - \overline{B}^0$ and $B^0 - \overline{B}^0$ oscillations in which the mesons turn into their antimatter partners many times during their lifetime, the $D^0 - \overline{D}^0$ oscillations are very slow, over one hundred times the average lifetime (see <u>7 November 2012</u> news for introduction). LHCb has now updated this result using the full 2011 and 2012 data set of 3 fb⁻¹. The new result is 2.5 times more precise. The values parameterizing the oscillations, the so-called mixing parameters y' and x'², are shown above.

By now, CP violation, differences in the behaviour of matter and antimatter, has been observed in all oscillating neutral-meson (K^0 , B^0 , B^0_s) systems apart from the charm system. First evidence for charm CP violation (see <u>14 November 2011</u> news) has not been unambiguously confirmed to date (see <u>12 March 2013</u> news). The D⁰ mesons are the only mesons containing up-type quarks which undergo matter anti-matter oscillations (called also mixing) and therefore provide unique access to effects from physics beyond the Standard Model.

As part of the new analysis, LHCb has investigated whether there is a CP violating contribution to the oscillations, in contrast to the Standard Model expectation. This is done by investigating whether the oscillation parameters for mesons produced as D^0 and \overline{D}^0 differ. Studying the D^0 and \overline{D}^0 decays separately shows no evidence for CP violation and provides the most stringent bounds on the parameters (A_D and |q/p| for experts) describing this violation from a single experiment.



(2) LHCb physicists measured the asymmetry A_Γ of the inverse of effective lifetimes in decays of D⁰ and \overline{D}^0 mesons to the K⁻ K⁺ and π⁻π⁺ final states. The measured values of the parameter A_Γ shown above represent the world's best measurements of this quantity, and are the first searches for CP violation in charm oscillations with sensitivity better than 10⁻³. They do not indicate CP violation, and show no difference in A_Γ between the two final states.

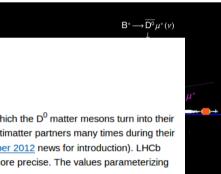
The results of other experiments combined by the <u>Heavy Flavor Averaging</u> <u>Group</u> indicated a hint for possible non-zero values of the CP violation parameters (|q/p| and ϕ for experts). Both LHCb results presented above do not support this indication as seen in the image. The size of the contour with the new LHCb results is about a factor of two smaller in each of |q/p| and ϕ . They provide very stringent limits on the underlying parameters, thus

12 March 2013: Improved search for CP violation in charm decays.

$[\Delta A_{CP} = (-0.34 \pm 0.15 \pm 0.10)\%, \text{ pion tagged }] \\ [\Delta A_{CP} = (+0.49 \pm 0.30 \pm 0.14)\%, \text{ muon tagged }]$

The LHCb Collaboration presented today at the <u>Rencontres de Moriond QCD</u>. La Thuile, Italy, results of an improved search for the difference between properties of matter and antimatter, CP violation, in charm decays, see <u>14 November 2011</u> news for introduction. The difference (Δ) of CP asymmetry (A_{CP}) between the decay rates of D (matter) and \overline{D} (antimatter) mesons into <u>K^*K</u>⁻ pairs and into <u>m^*n</u>⁻ pairs was measured. The results presented today profited from three improvements to the previous analysis: the full 1.0 fb⁻¹ data sample collected in 2011 was used instead of 0.6 fb⁻¹, the analysis technique was improved and also in addition another independent method was used to select matter D and antimatter D particle decays.

In the Standard Model CP violation was expected to be very small in the charm sector, whereas new physics effects could generate enhancements. Therefore the <u>14 November 2011</u> announcement by the LHCb Collaboration of <u>3.50</u> evidence of CP violation in charm sector, $\Delta A_{CP} = (-0.82 \pm 0.21 \pm 0.11)_{\%}$, triggered intensive theoretical activity with conclusions that some special Standard Model effects could generate CP violation effects even as big as about <u>1%</u>. This interesting LHCb result was later confirmed by the CDF and Belle collaborations. The new improved LHCb result presented today, $\Delta A_{CP} = (-0.34 \pm 0.15 \pm 0.10)_{\%}$, is more precise thanks to the larger data sample and several improvements resulting in better background suppression by a factor of 2.5. The central value is, however, closer to zero than in the previous measurement, which it supersedes.



In the measurement presented above the D (matter) and \overline{D} (antimatter) mesons were selected using the D* meson decays, $D^{*+(-)} \rightarrow \pi^{+(-)}D(\overline{D})$, which means that the presence of π^{+} in the decay identified matter D meson production while π^{-} accompanied antimatter \overline{D} production. LHCb physicists presented today also results of a second independent analysis in which the D and \overline{D} mesons were selected using so called semileptonic beauty B decays, for example $\mathsf{B}^{+(\text{-})} \to \mu^{+(\text{-})} \nu \overline{\mathsf{D}}(\mathsf{D}).$ In the second analysis, the positive charge of μ^+ identified the \overline{D} meson, while the negative one, μ , the D production. The image at the left hand side shows a selected event. A zoom around the pp interaction point shows a B⁺ meson decay point located at the distance of 17 mm from the pp collision point and the \overline{D} meson decay place still 9 mm further away. The second analysis also measures a value that is consistent with zero: $\Delta A_{CP} = (+0.49 \pm 0.30 \pm 0.14)\%$. A combination of the two LHCb results gives $\Delta A_{CP} = (-0.15 \pm 0.16)\%$.

Ceard, CP3, Petnica 2014